

From Inherent Optical Properties to Biogeochemical Properties

Emmanuel Boss, U. of Maine

- Forward and inverse problems.
- The basis for the relationship (theoretical, empirical, hybrid).
- Concept of a 'proxy'.
- Bulk vs. single particle property (FCM).
- Supportive lab studies (controlled compared to ocean).
- Extensive vs. intensive properties.
- Some intensive proxies involve ratio of proxies.
- Uncertainties...

Remember: this is the major reason the field of Oceanography cares about optics!!!

1st order variability in optical properties is due to concentration (optical parameter are additive ← Beer-Lambert-Bouger law).

What is the range of changes in concentration?

What else affects optical properties (2nd order variability)?

Composition (index of refraction)

Size (what is size?)

Shape (e.g. axis ratio, micro)

Internal structure (e.g. cell wall, organelle)

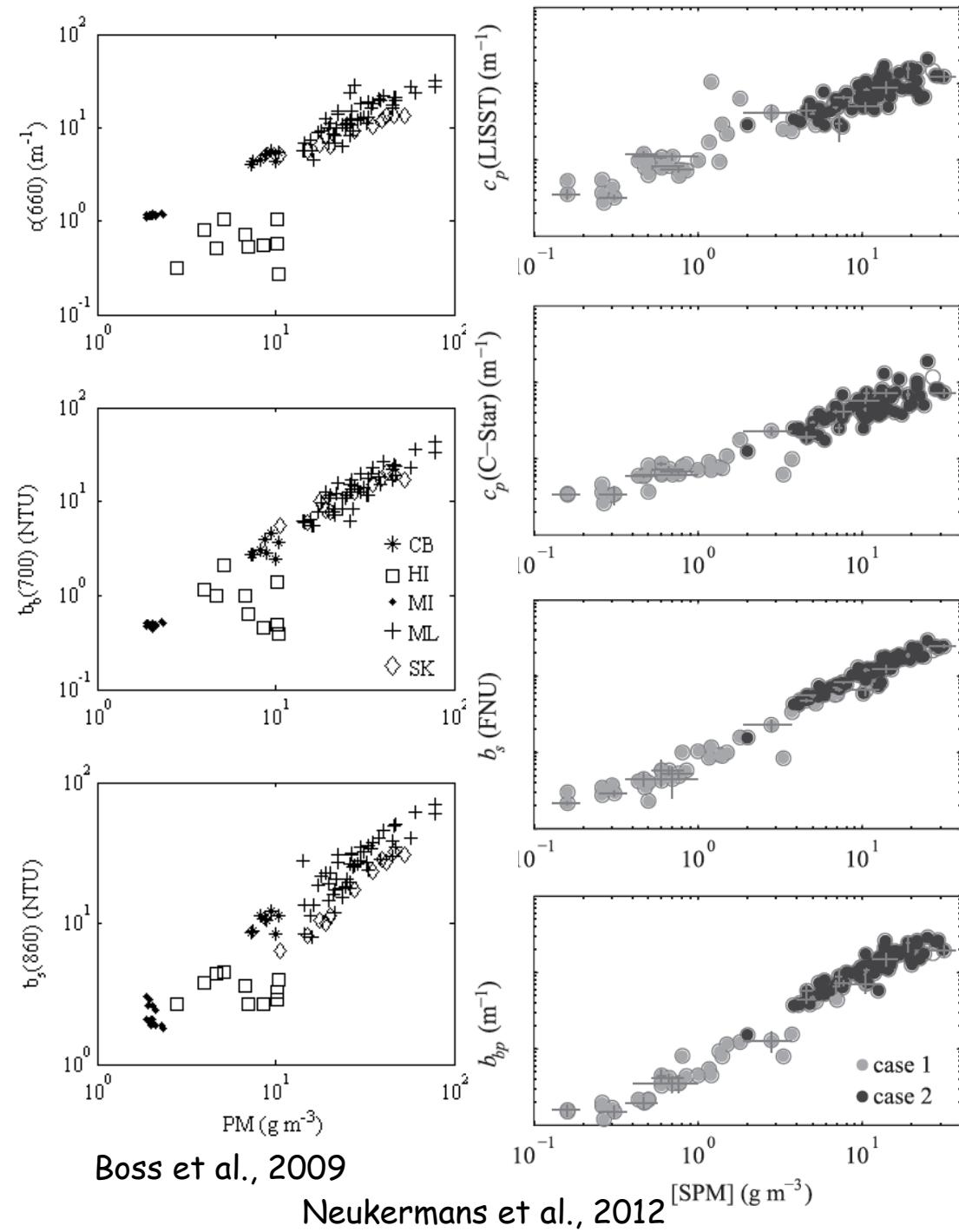
Packing (fluid fraction in aggregate)

Proxies for particulate mass (extensive)

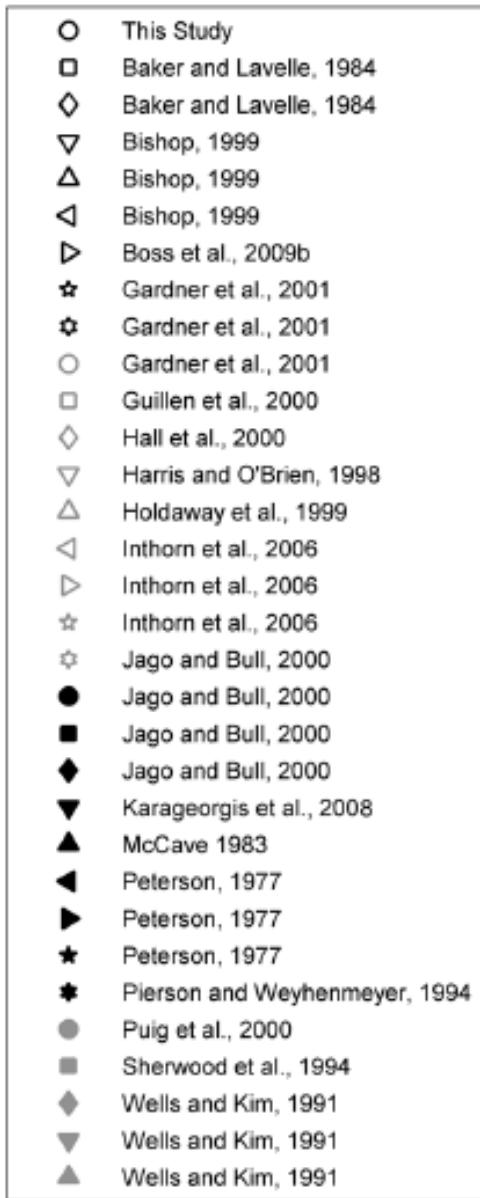
- Many comparisons, starting in the 70s.
- Used to study sediments (signal to noise).
- Moved to open ocean as calibration/stability improved.

Which regression type should one use for a proxy?

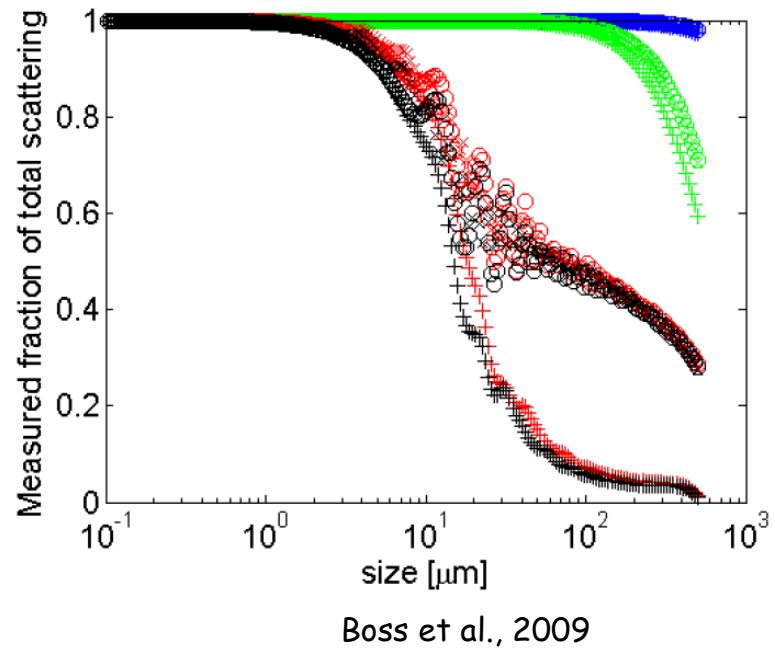
Optics is a standard method to measure turbidity, a primary determinant of water quality (e.g. ISO-7027).



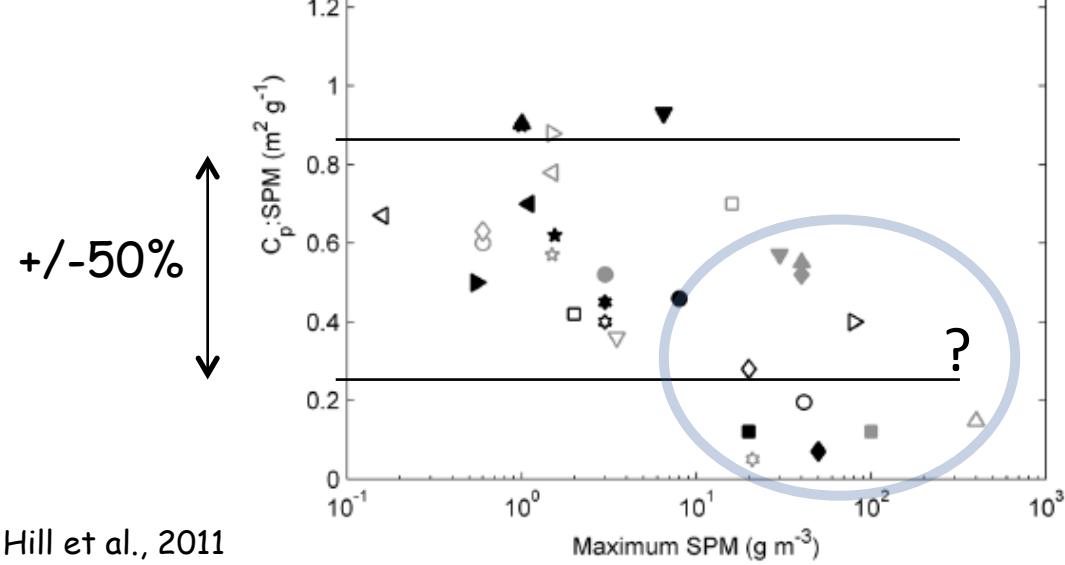
$c_p(660)$ vs. mass:



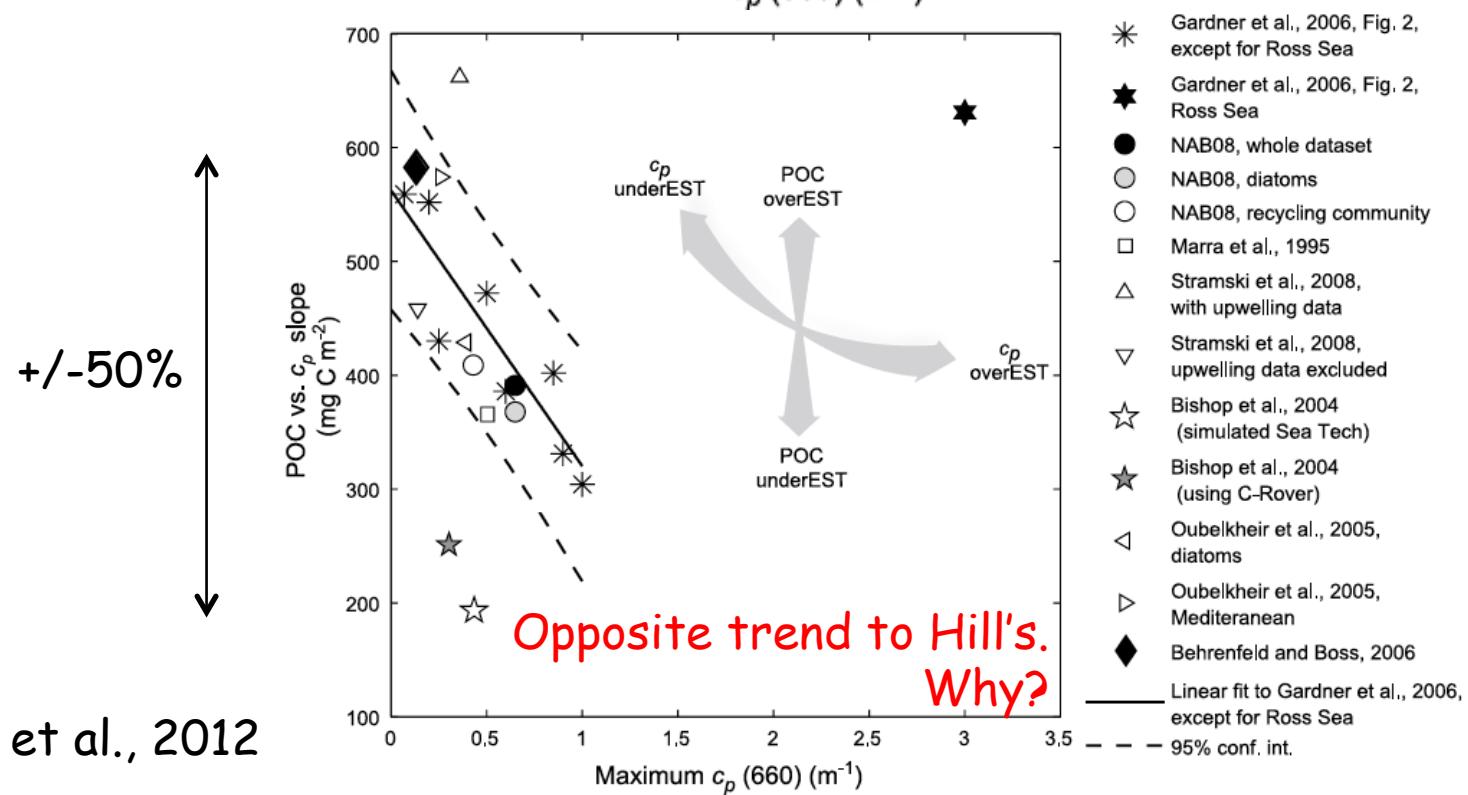
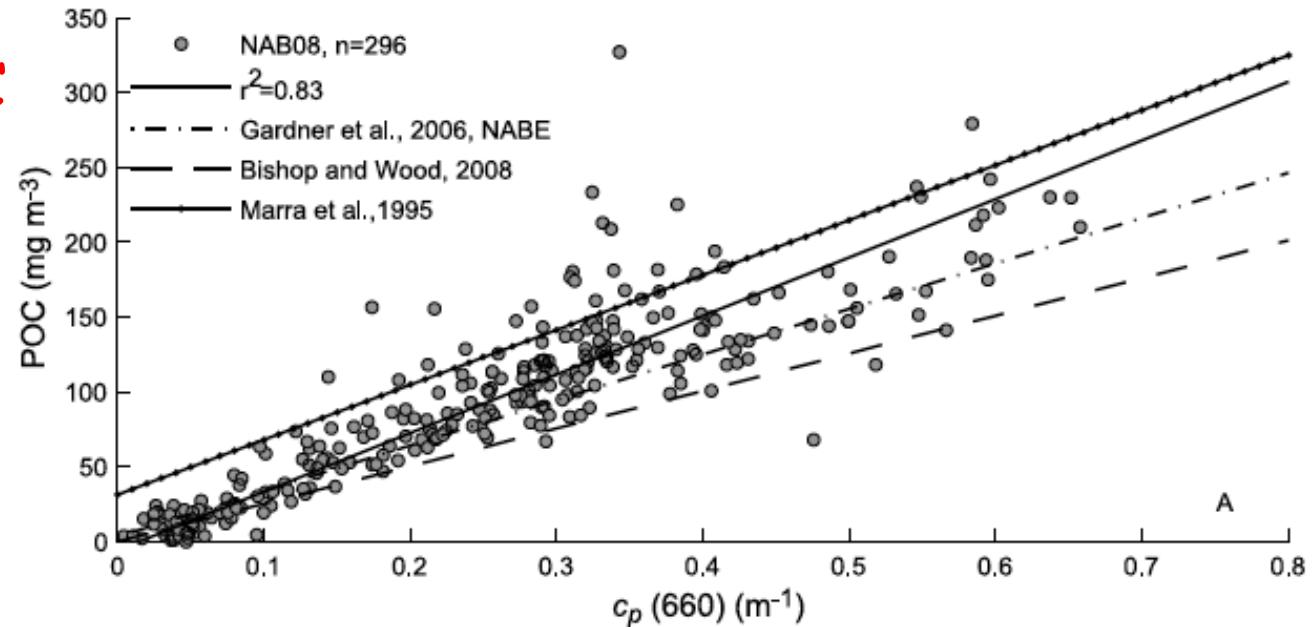
Theory:



Observations:

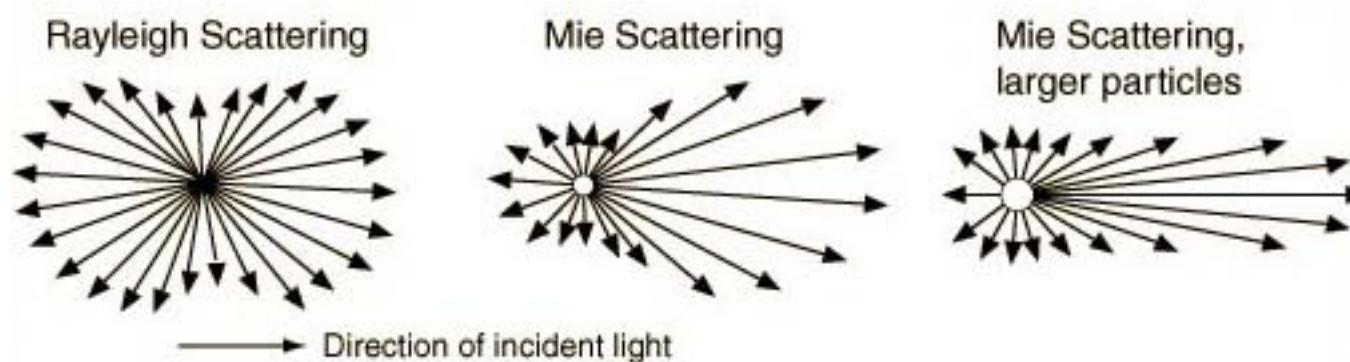


$c_p(660)$ vs. POC
(extensive)



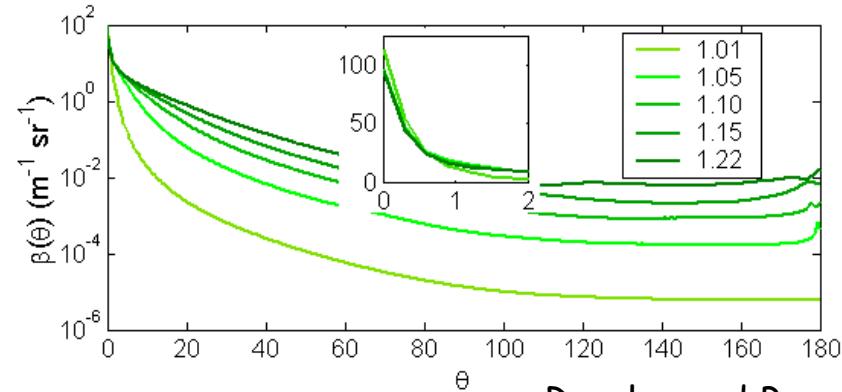
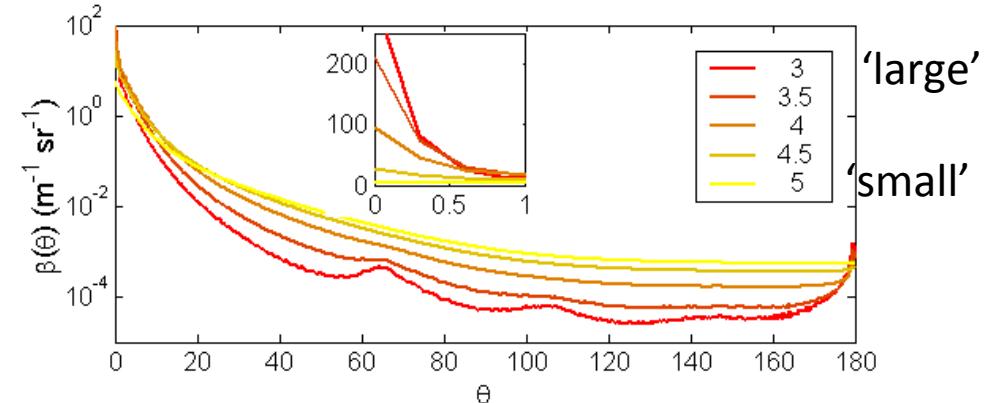
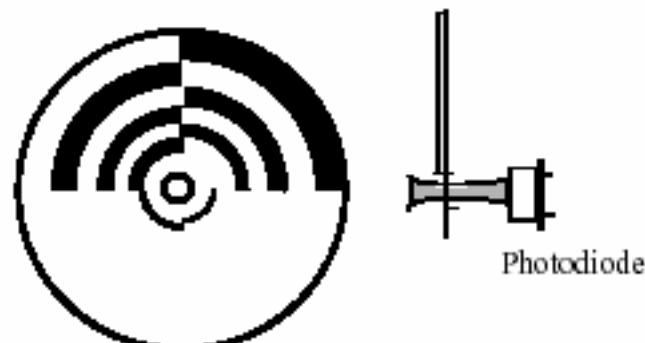
Cetinic et al., 2012

Angular dependence of scattering on size



Near forward scattering: Strong dependence on size, less on n .

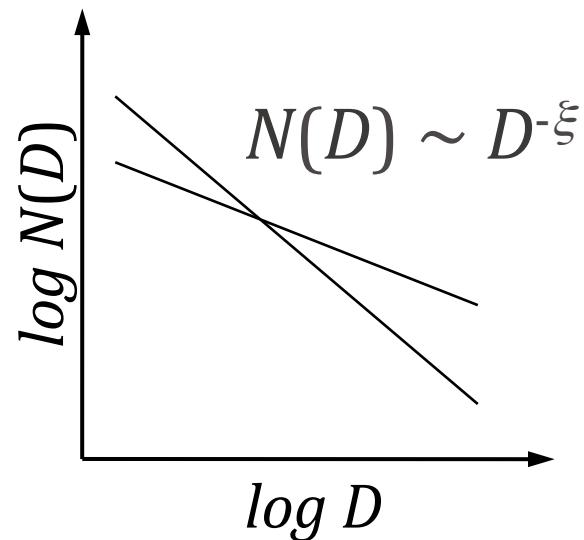
LISST detector:



Roesler and Boss, 2008

Spectral c_p

(1) Assuming a power-law particle size distribution (PSD)



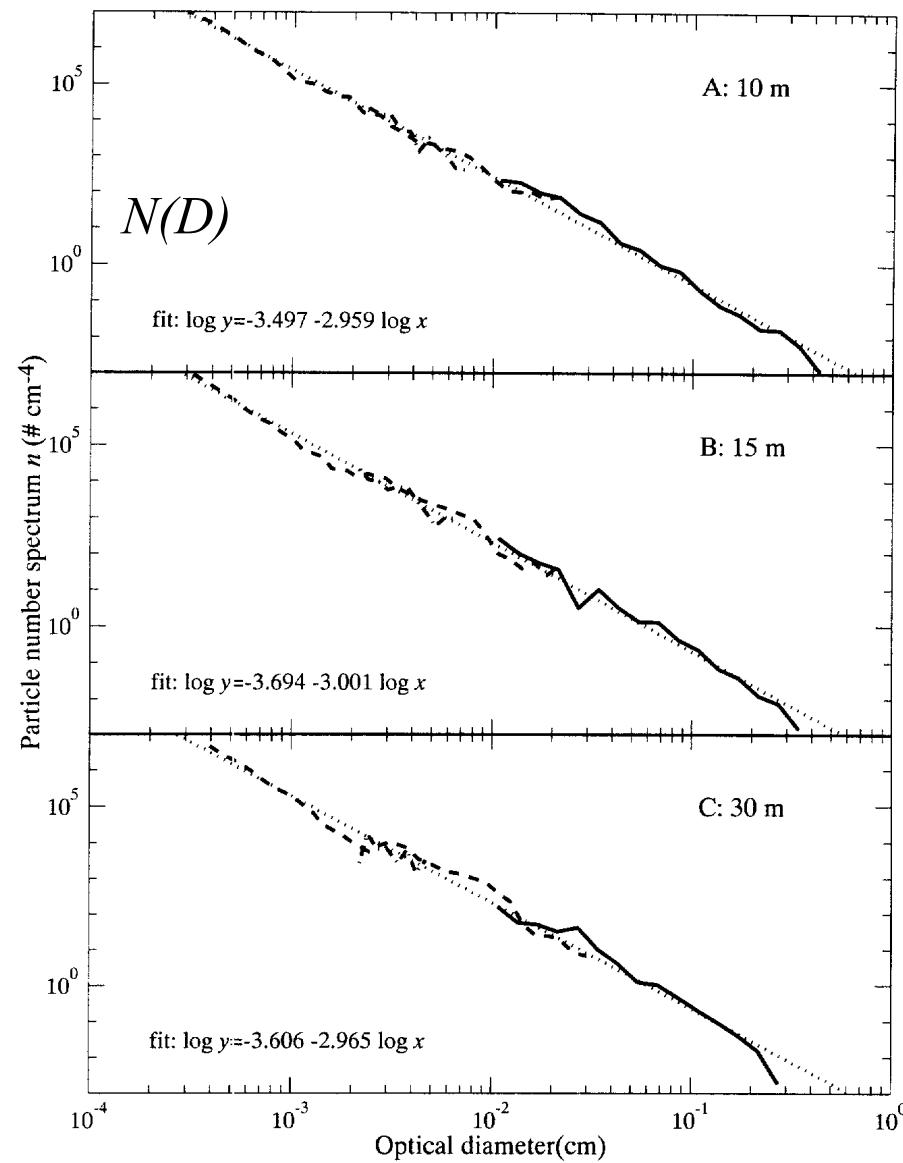
→ $c_p(\lambda)$ is described well as a power law function of wavelength (λ)

$$c_p(\lambda) \sim \lambda^{-\gamma}$$
$$\gamma \approx \xi - 3$$

→ Flatter beam attenuation spectra (small γ) implies flatter particle size distribution (small ξ)

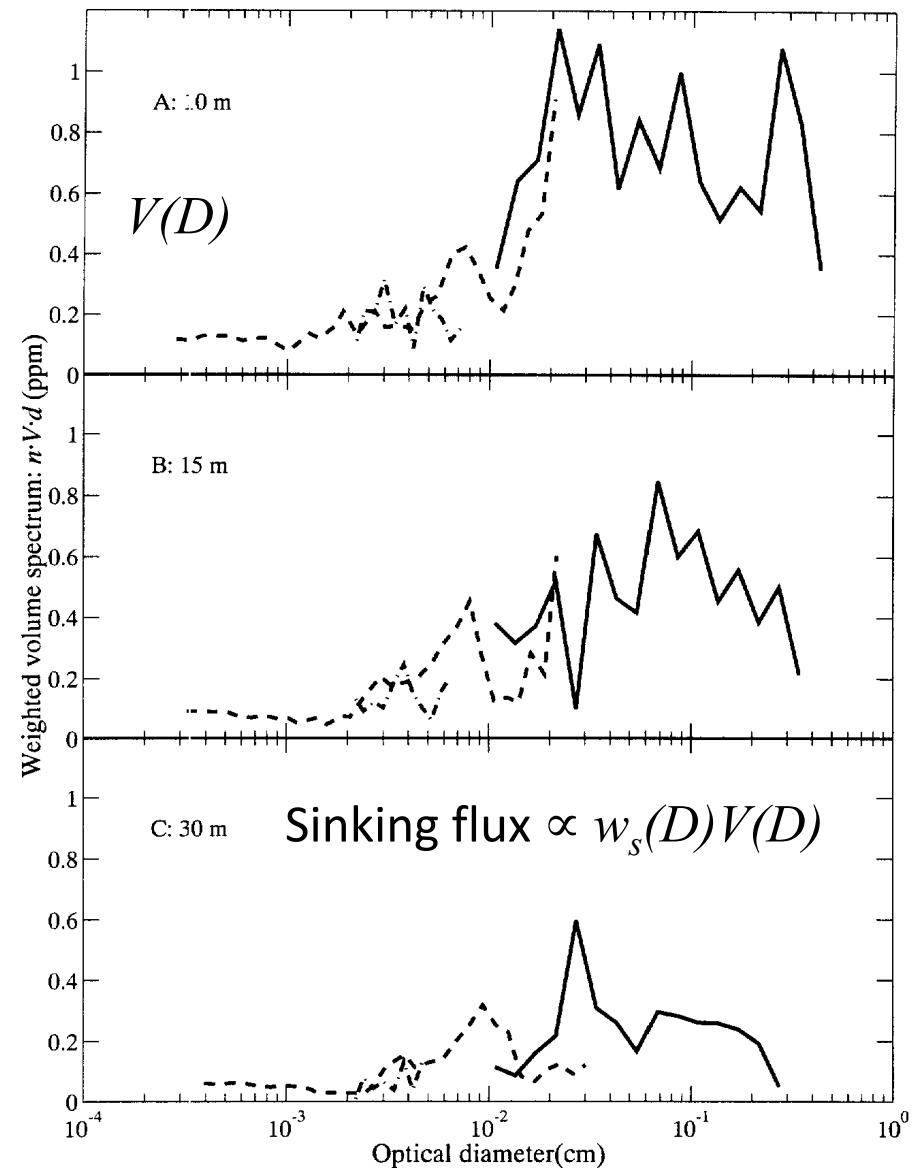
(2) Assuming spherical non-absorbing particles

Particle size spectra between 1 μm and 1 cm at Monterey Bay determined using multiple instruments

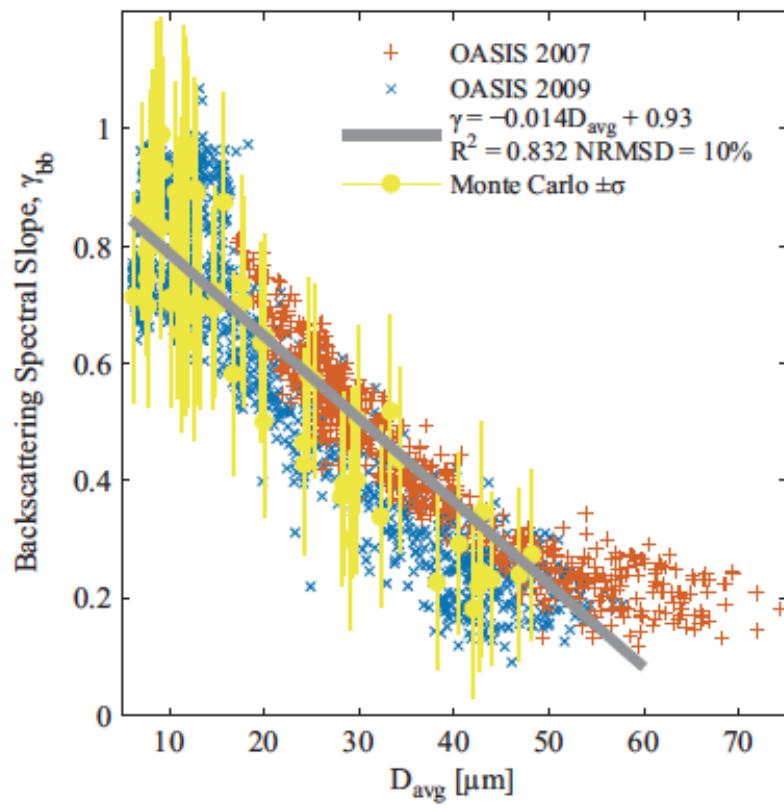


Jackson et al., 1997

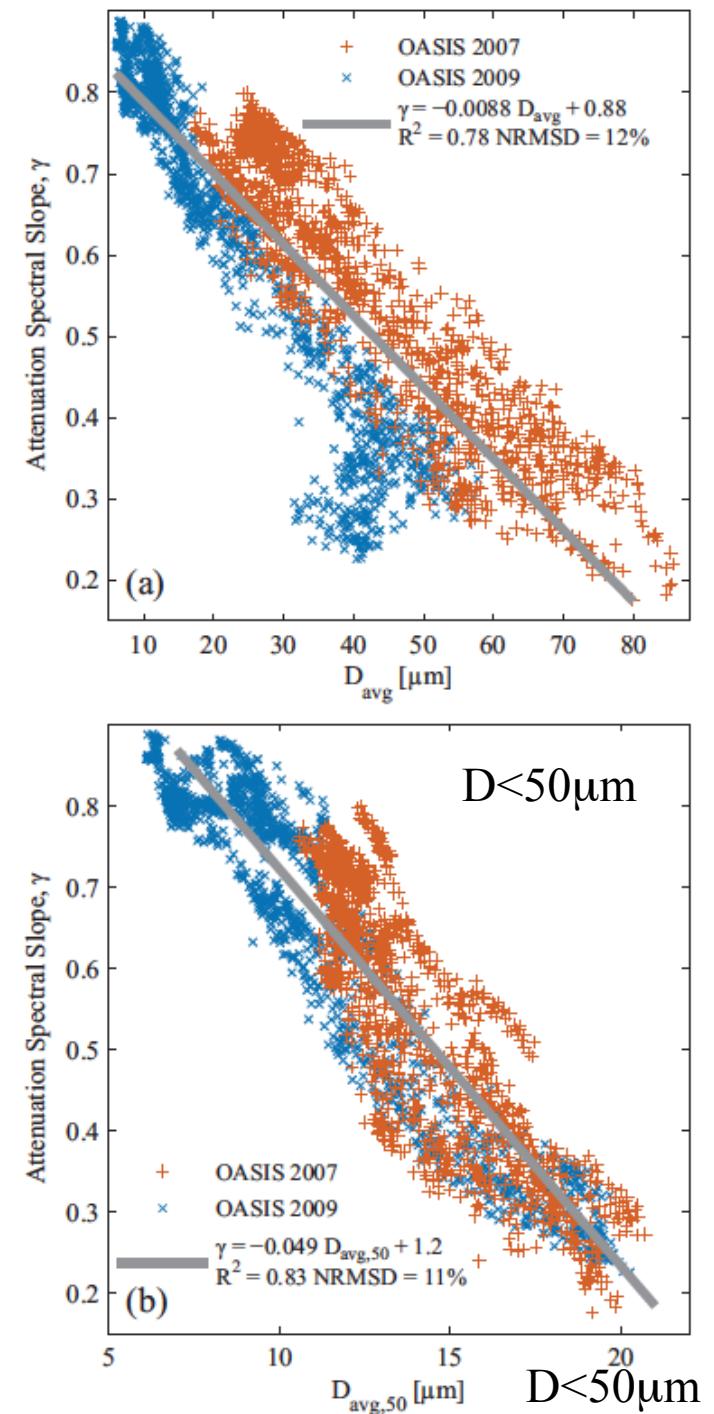
Focus is different



Closure: LISST vs. IOP spectra Field data, MVCO



Slade and Boss, 2015



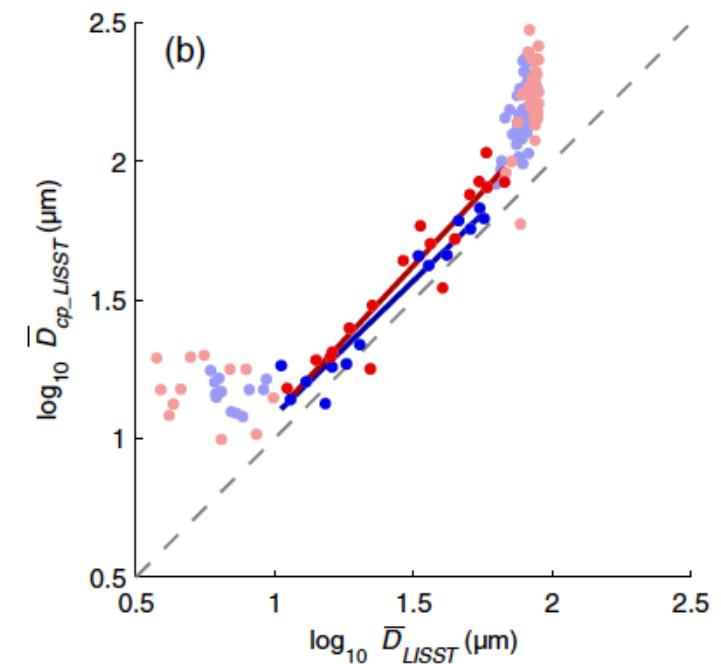
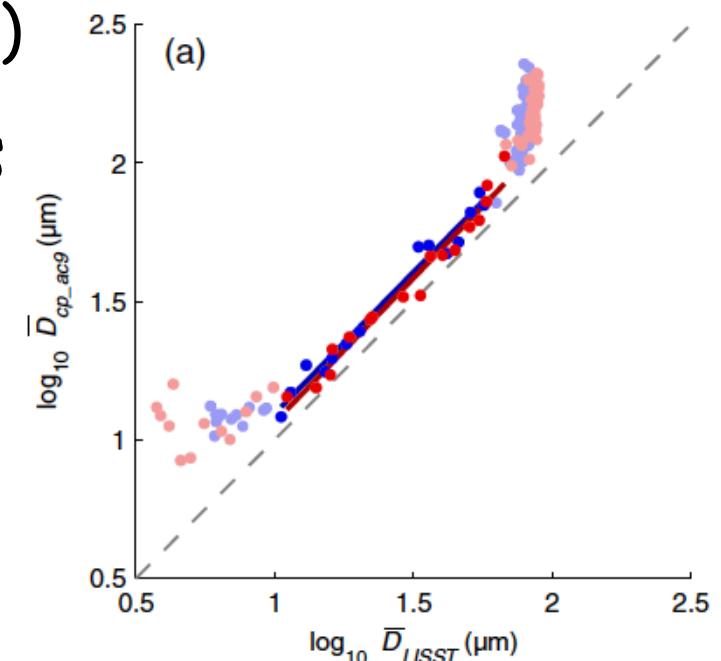
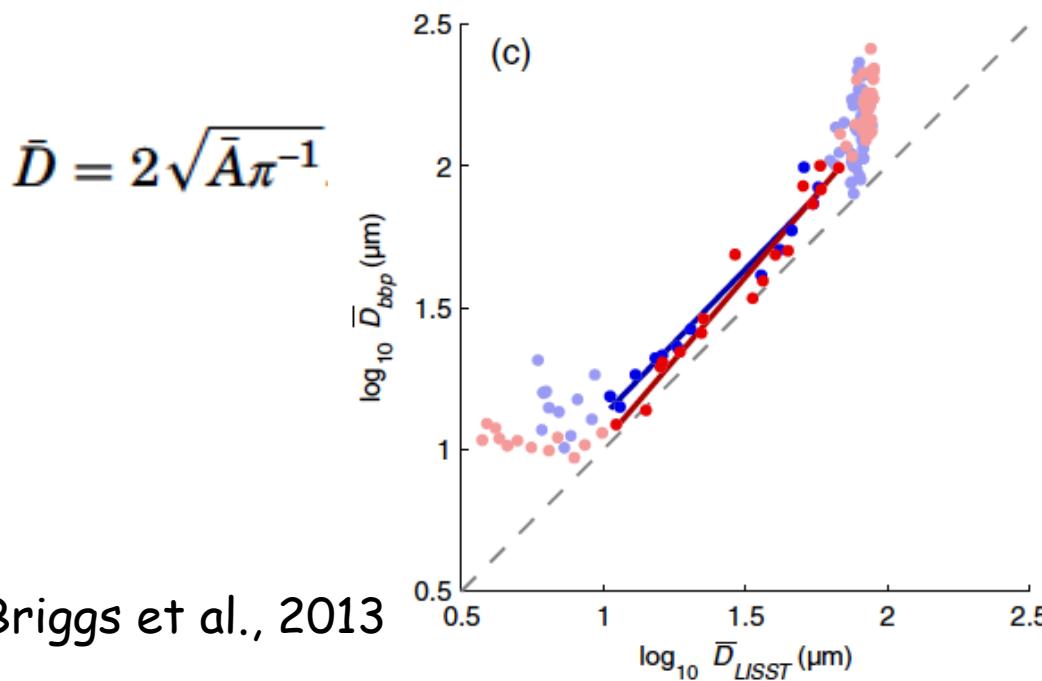
Method of fluctuation (Shifrin, 1988)

Closure: LISST vs. fluctuations
Lab aggregation exp.

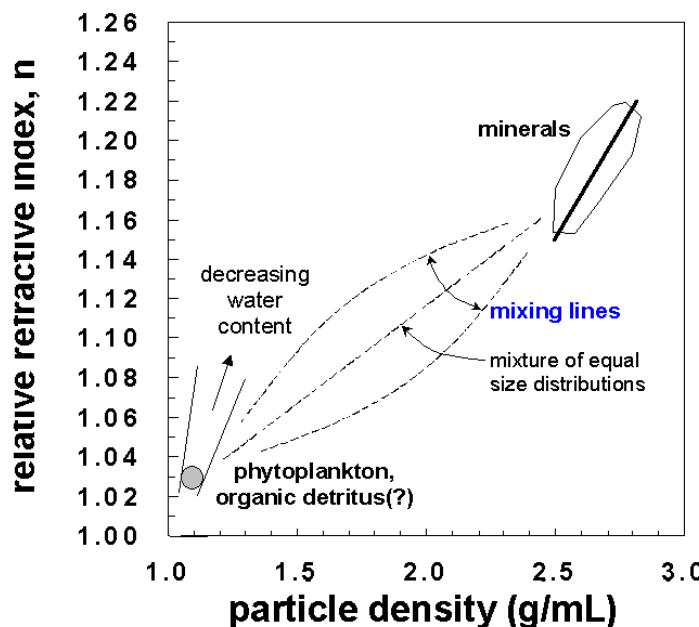
Sample volume

$$\bar{A}_{cp} = \frac{\text{Var}[c_p(t)]}{E[c_p(t)]} V \frac{1}{Q_c \alpha(\tau)}$$

Measurement time



Composition - index of refraction (an intensive parameter)



Zaneveld et al., 2002, OOXVI.
 Compiled from:
 Aas (1983)
 Carder et al. (1972)
 Carder et al. (1974)

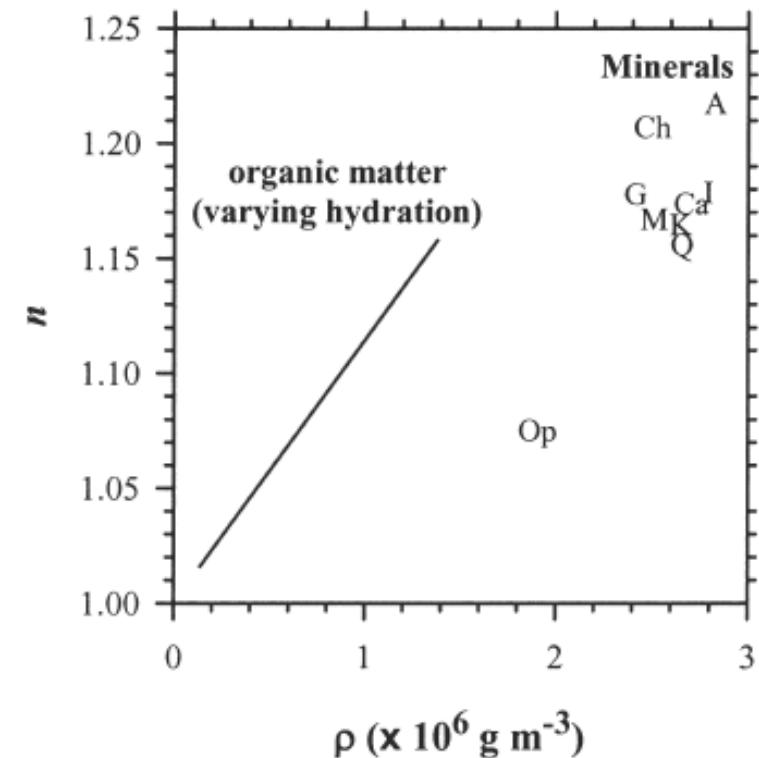
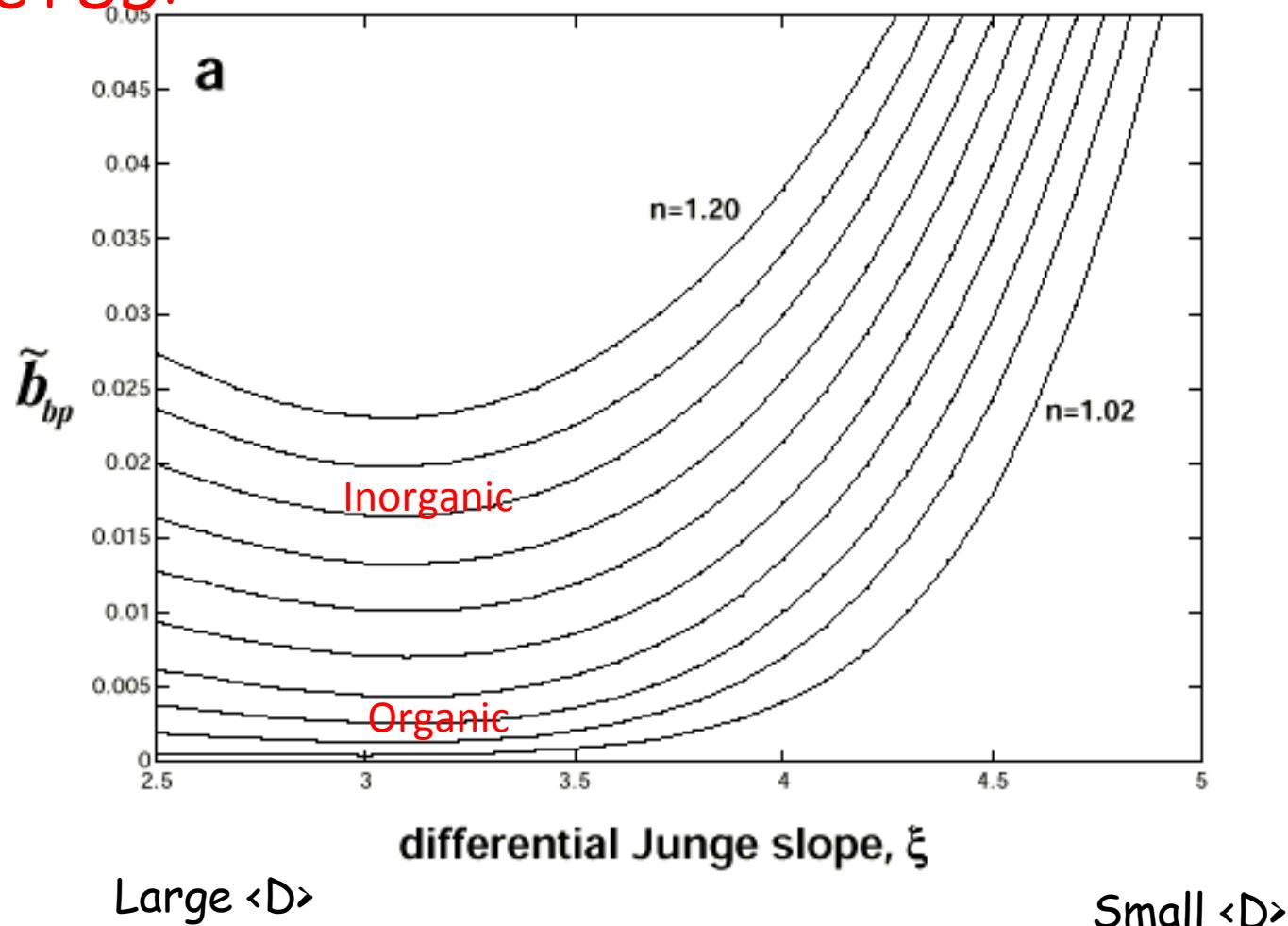


Fig. 7. The index of refraction relative to seawater for various minerals as a function of particle density. The n and ρ values (also listed in Table 6) are from Lide (2001). The n values are the arithmetic average of the values given for the two or three structure coordinate axes. The plotted minerals are aragonite (A), calcite (Ca), chlorite (Ch), gibbsite (G), illite (I), kaolinite (K), montmorillonite (M), opal (Op), and quartz (Q). The theoretical relationship between n and ρ is also shown for organic matter (see text for details).

Babin et al., 2003

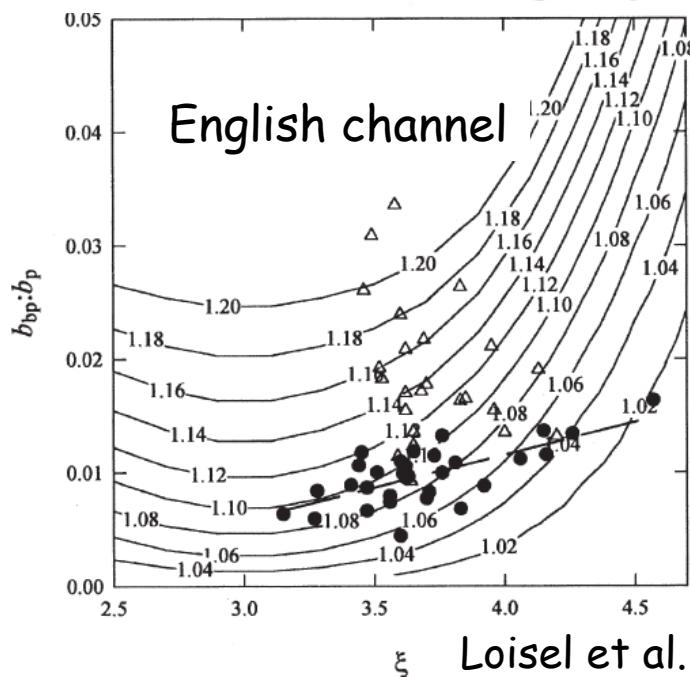
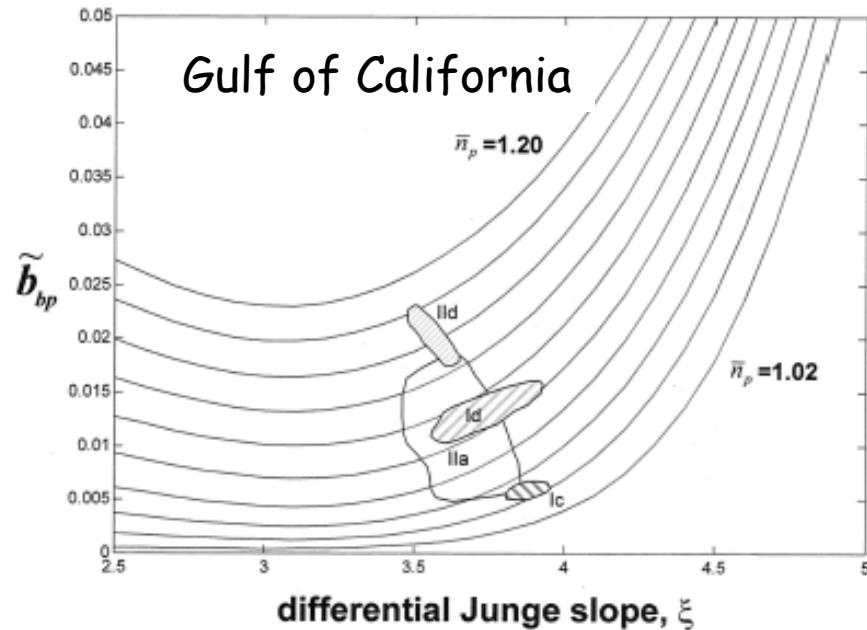
The making of a composition proxy:
Mie theory: the $b \sim b_{bp} = b_{bp}/b_p$ is very sensitive to n and less
so to the PSD:



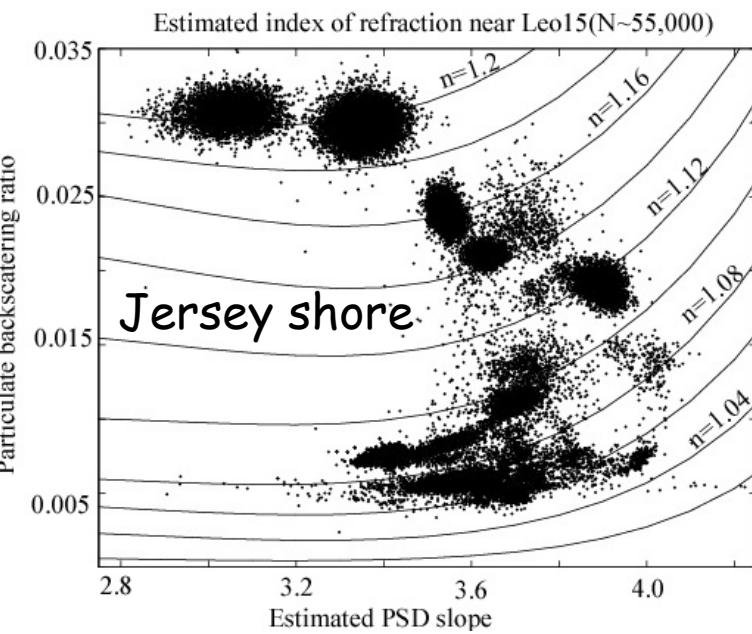
Twardowski et al., 2001

Twardowski et al., 2001:

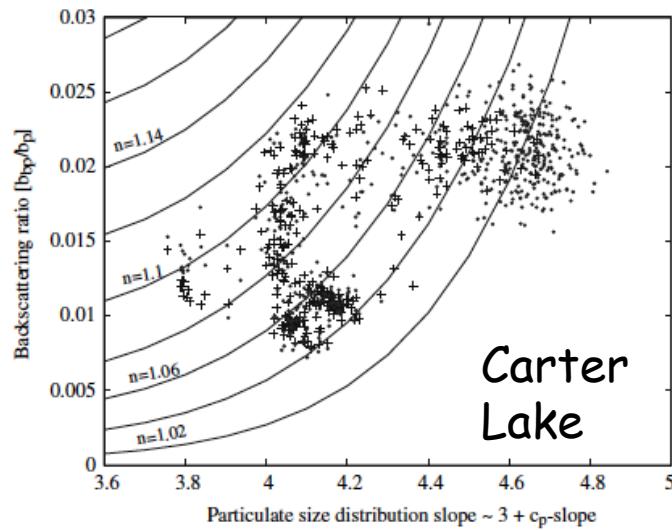
Observations



Loisel et al., 2004

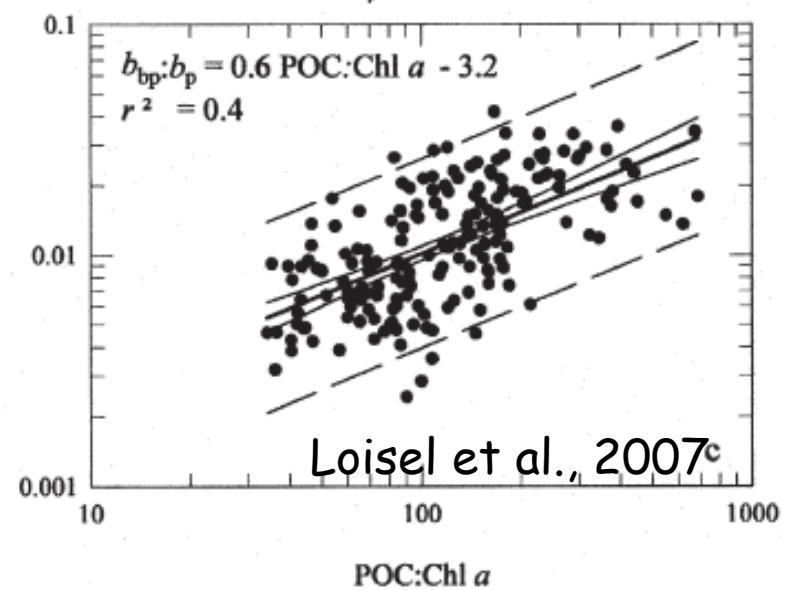
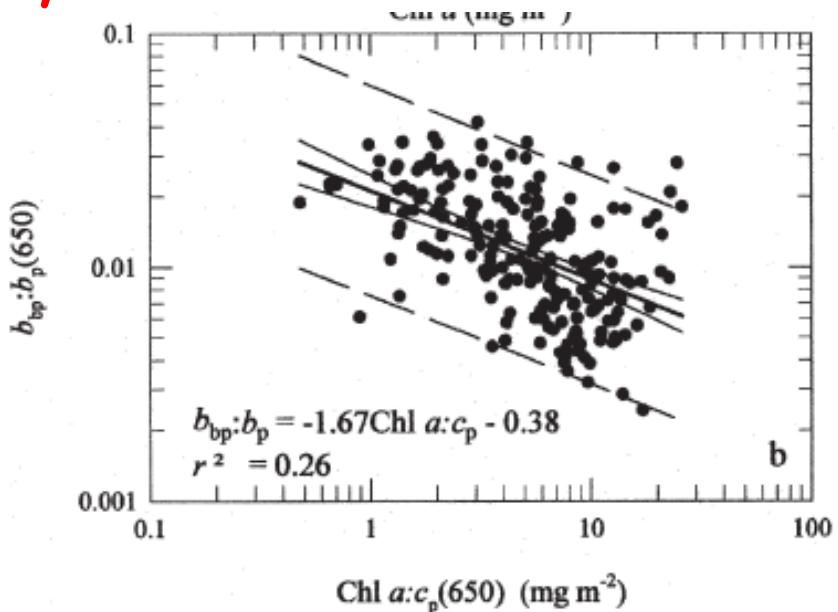
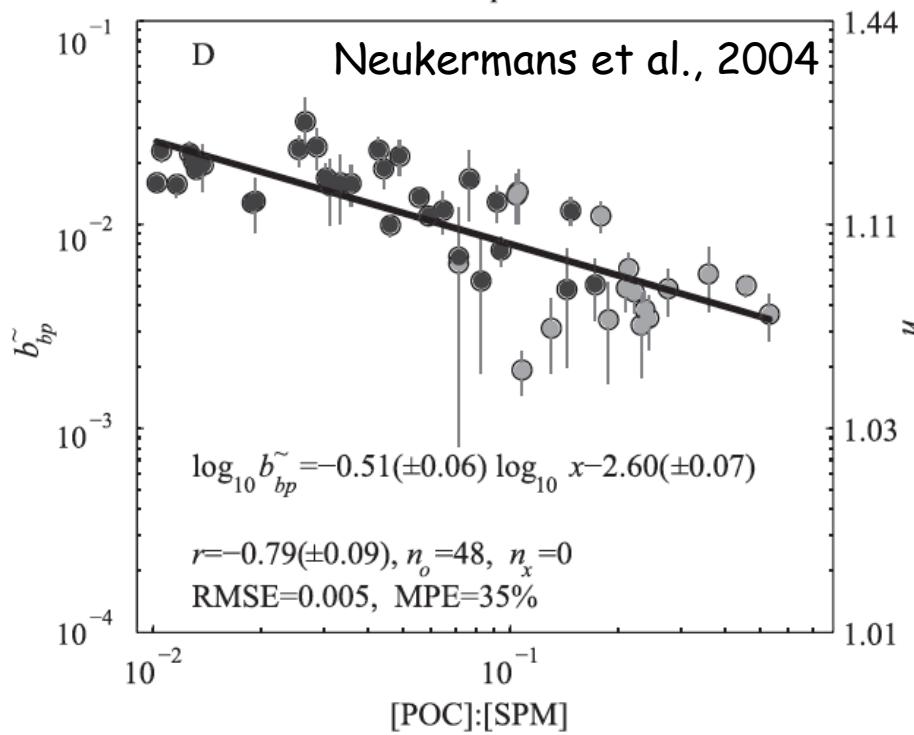
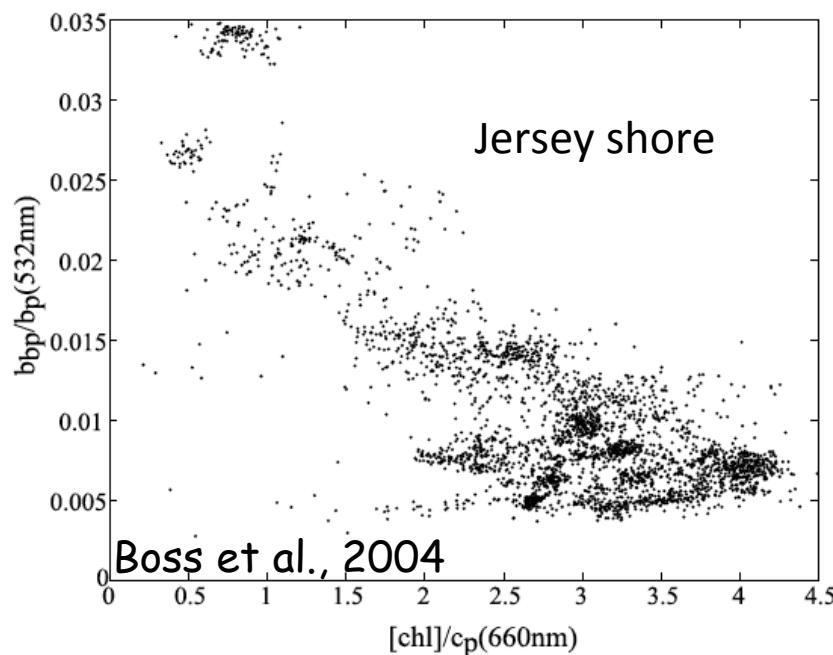


Boss et al., 2004



Boss et al., 2007

Proxy validation



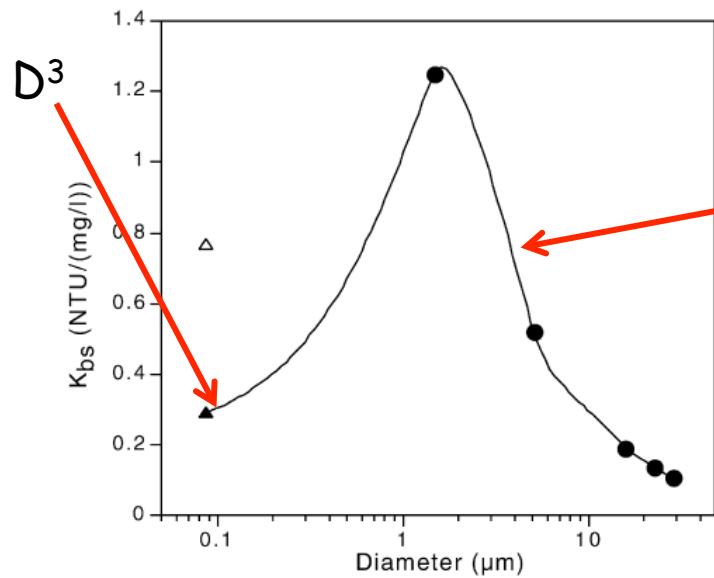
Varies from: phytoplankton → inorganic particles.

Mie theory tells us that the relationship between optical properties and mass is **composition** and **size** dependent:

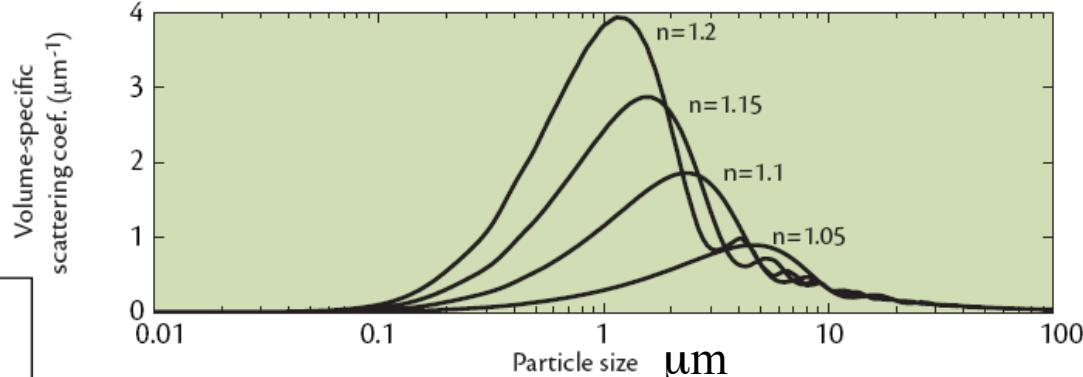
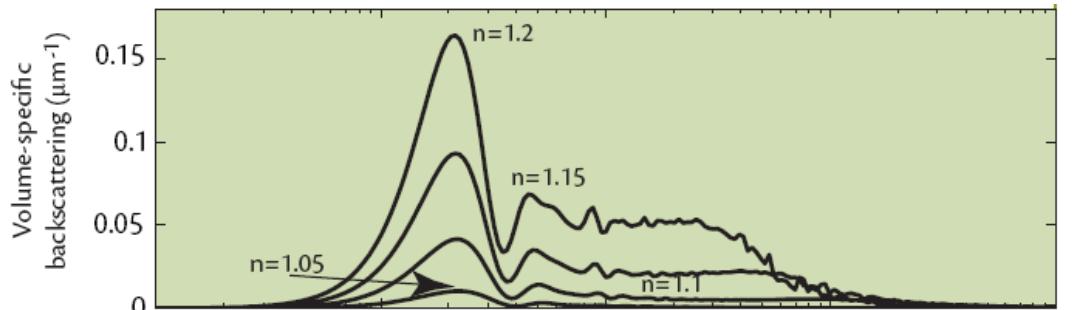
b_{bp}/Volume

Very different from:

b_p/Volume



Baker and Lavelle, 1984



Boss et al., 2004

b_{sp}/Mass

- All curves are 'resonant' curves

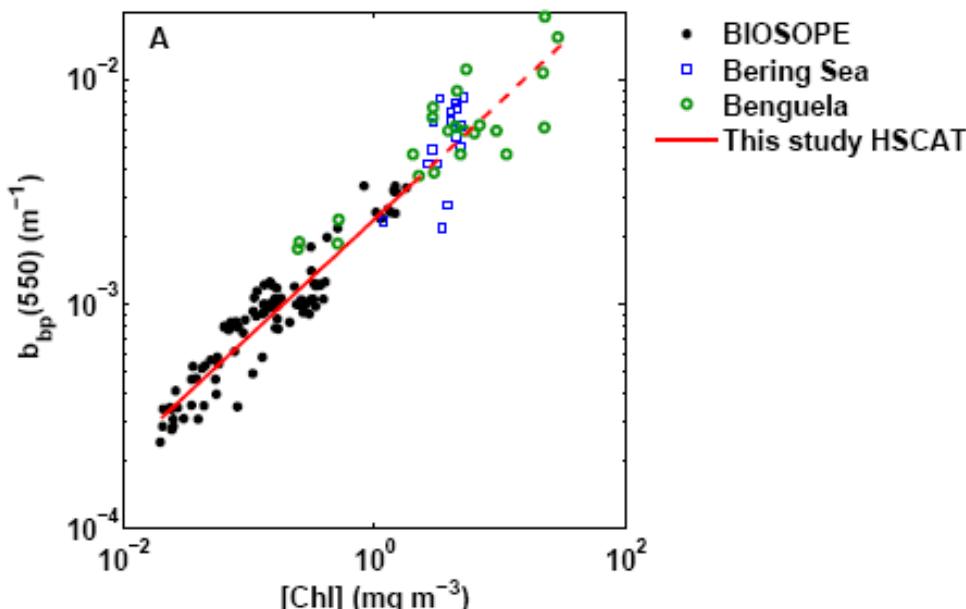
- Highest sensitivity for micron sized particles (c_p and b_s).

- Size of max response varies

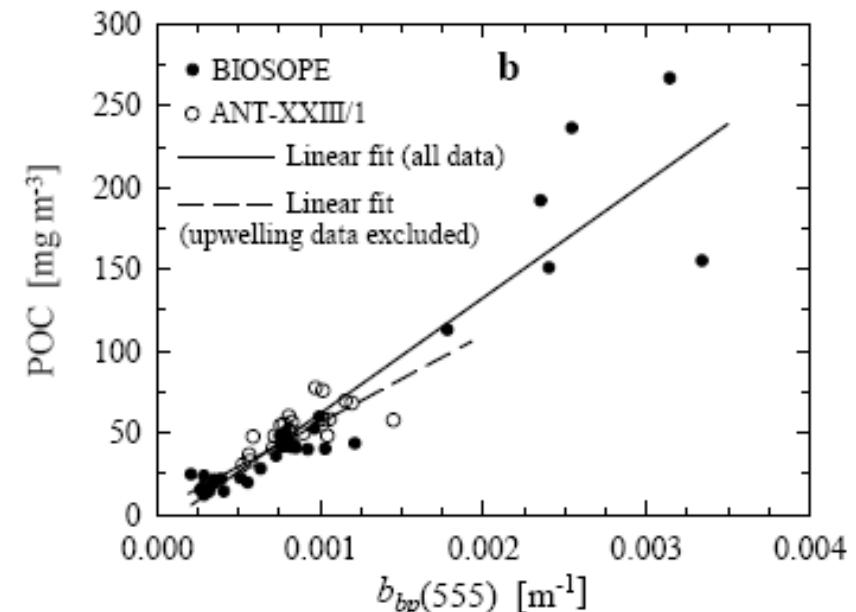
The b_b enigma (or paradox):

Based on Mie theory, backscattering should be dominated by inorganic particles and sub-micron particles (the least known).

Yet b_{bp} correlates well with [chl] and POC (>0.7mm):



Huot et al., 2008



Stramski et al., 2008

Possible explanation for the b_b enigma:

1. Mie results are correct. However, all particles in the open ocean co-vary, hence the tight relationship
← inconsistent with spectrum of b_{bp}/b_p .
2. Mie theory is not applicable. Organic particle actually backscatter more than we ascribe to them.
 - This last seems more consistent with size fractionated measurements (e.g. Dall'Olmo et al., 2009) and cultures (Whitmire et al., 2010, Poulin et al., 2018). Recent work supports modeling as coated sphere.

Backscattering by Nonspherical Particles: A Review of Methods and Suggested New Approaches

CRAIG F. BOHREN

Department of Meteorology, Pennsylvania State University, University Park

1991

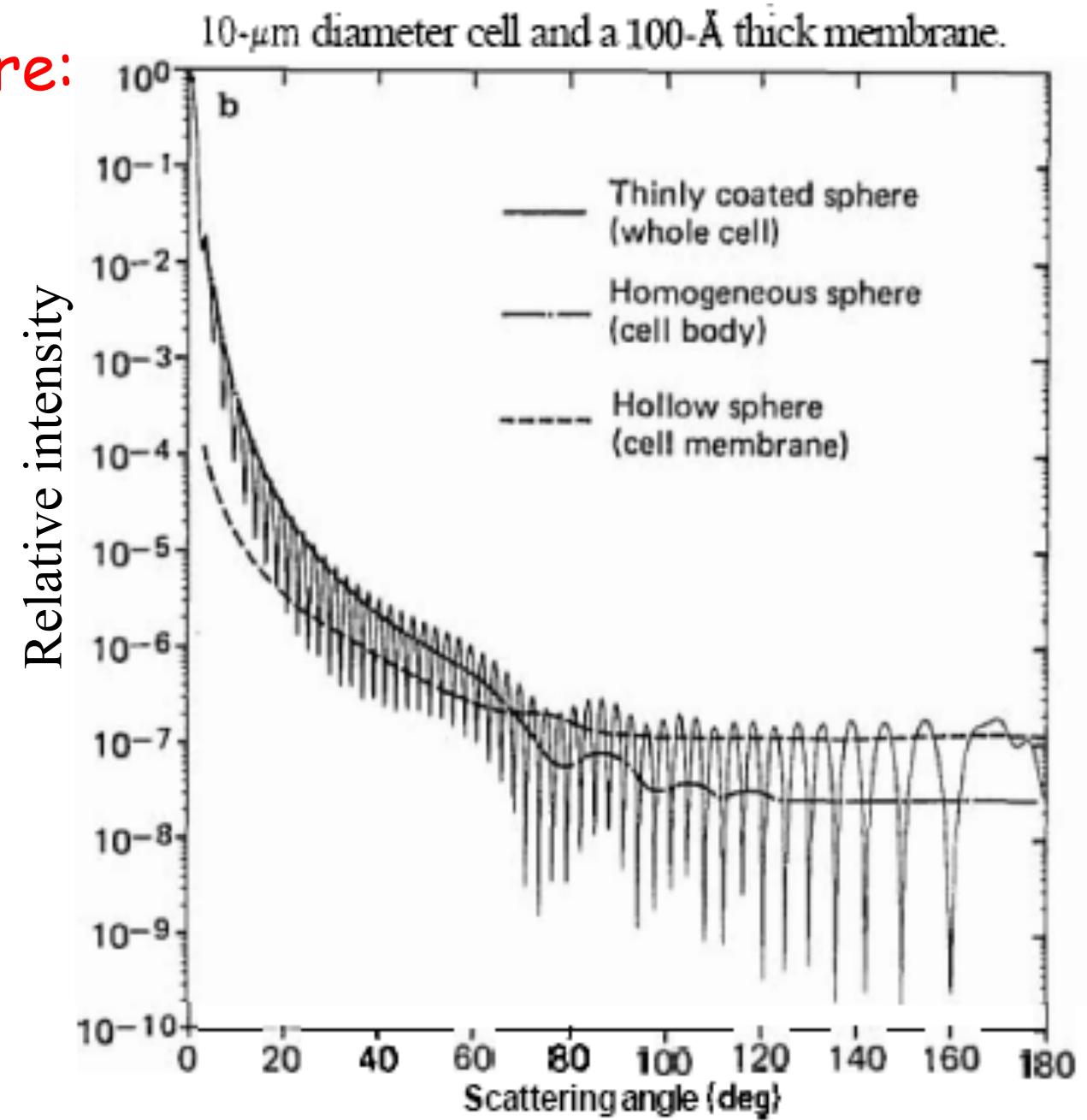
SHERMILA BRITO SINGHAM¹

Life Sciences Division, Los Alamos National Laboratory, Los Alamos, New Mexico

When criticized for using Mie theory where its applicability is dubious, modelers sometimes respond that although they know that Mie theory is inadequate, it is the only game in town. Better to do wrong calculations than to do none at all. Modelers have to model.

We suggest an alternative to modeling. It is called not modeling—not modeling, that is, until adequate methods are at hand.

Internal structure:



Backscattering dominated by membrane.

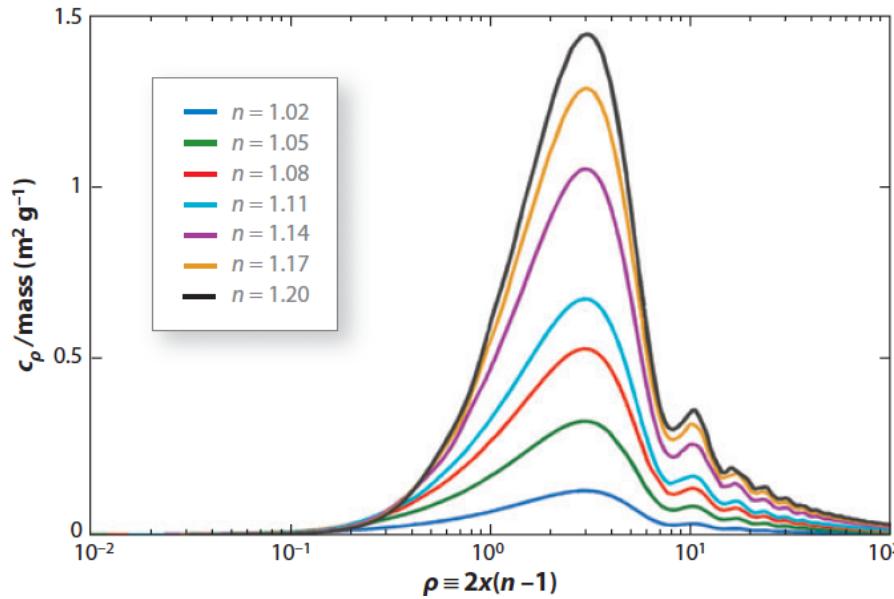
Meyer, 1979

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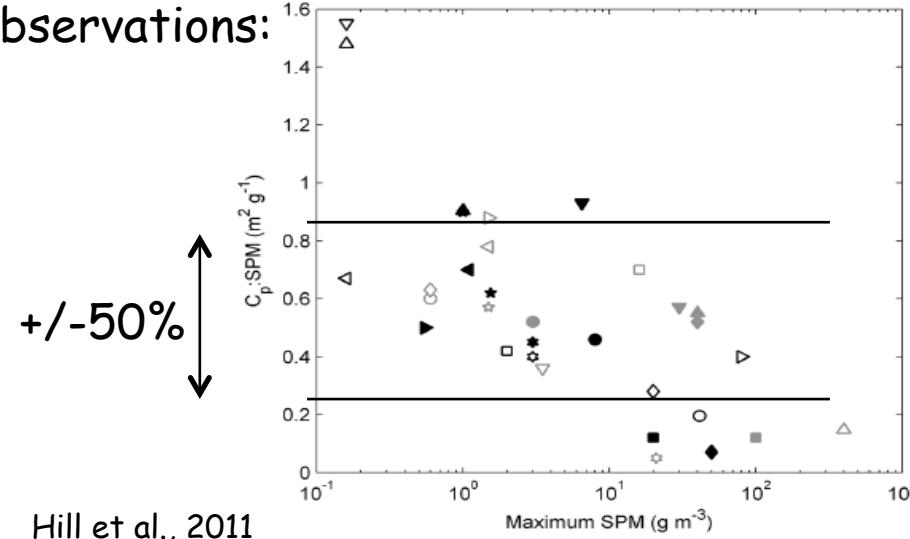
Back to c_p -SPM. Why so good?

If:



Observations:

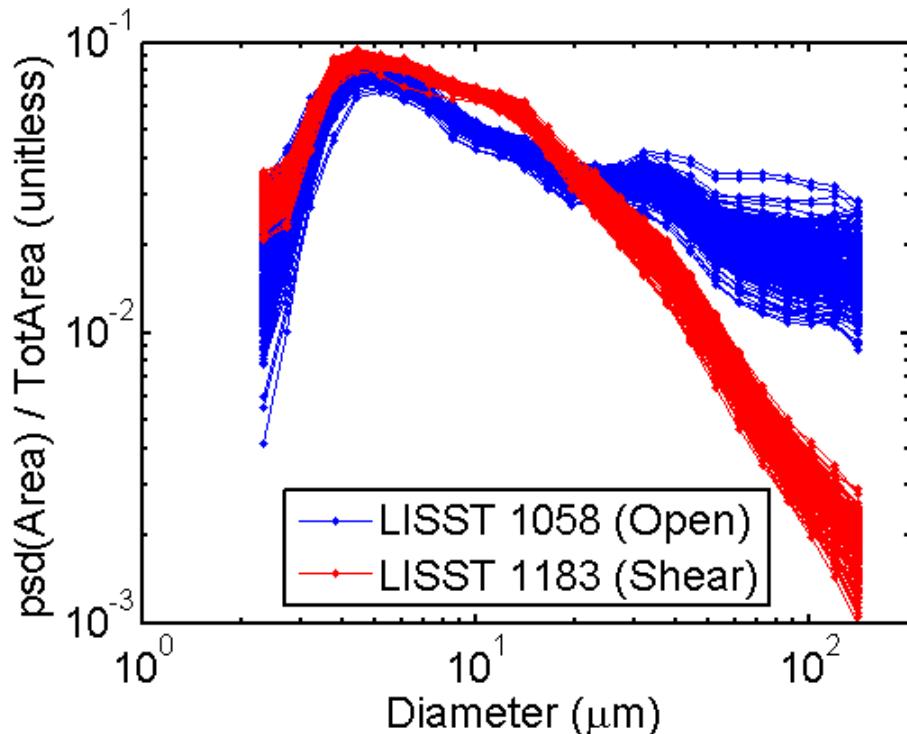
Why:



Hill et al., 2011

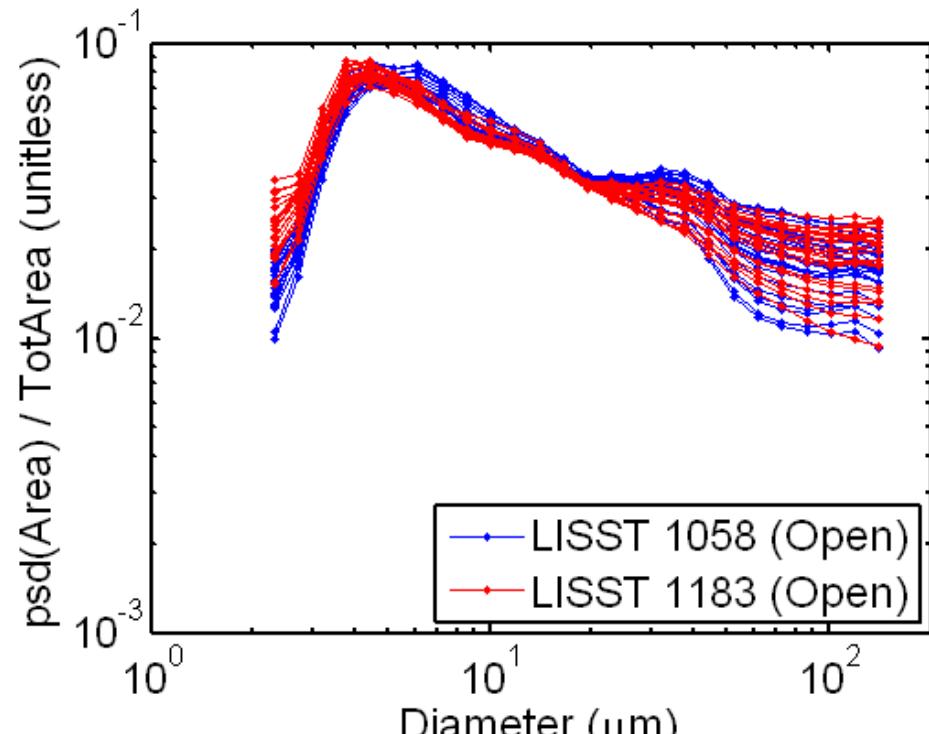
Hypothesis: aggregation reduces dependence of mass proxies on size.

Field manipulation:



Experiment

One LISST Pumped (red) and the other not (blue)

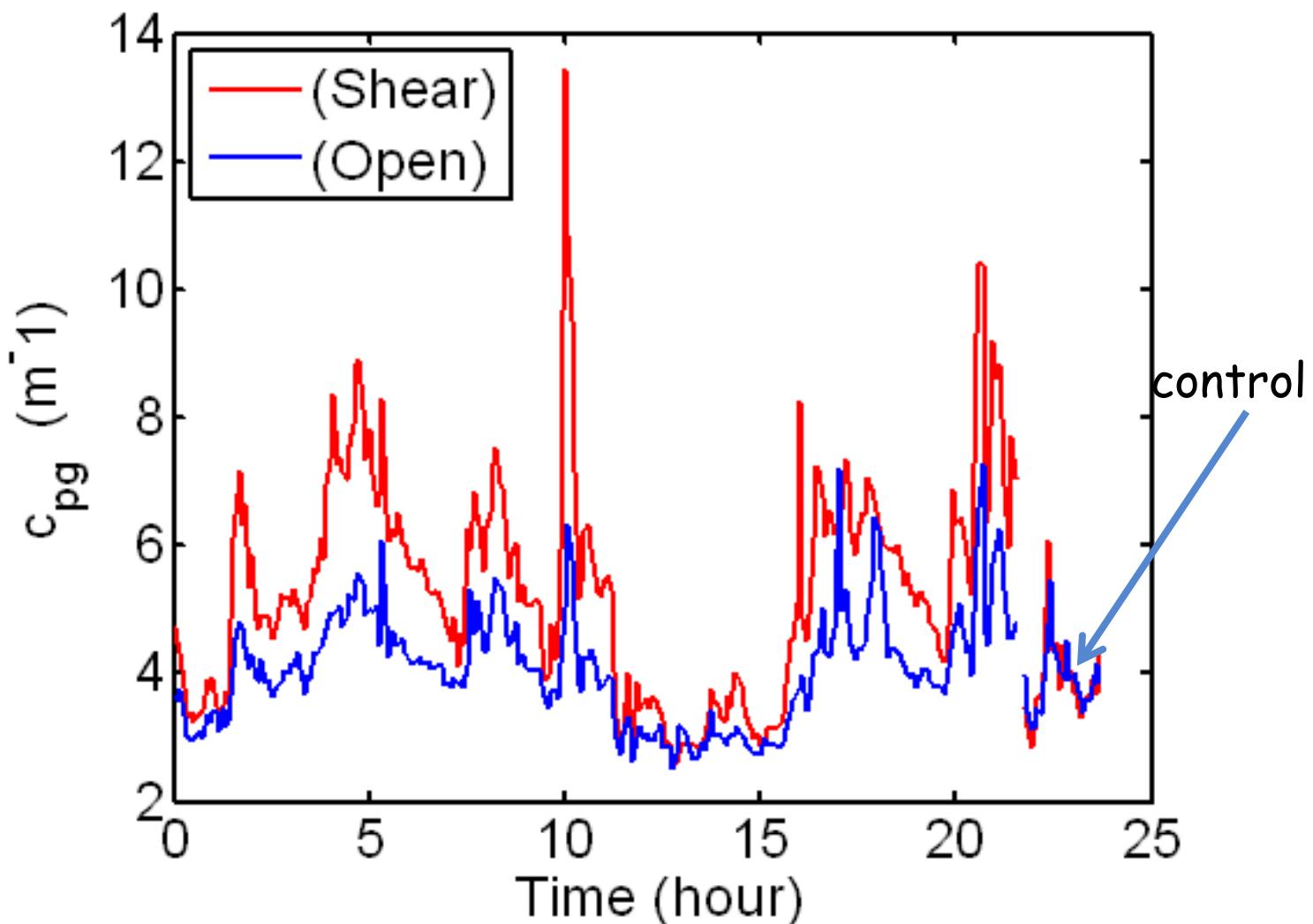


Control

both instruments deployed unpumped

Slade et al., 2011

Effect on beam attenuation (2m depth at DMC)



On average, observed beam attenuation increases by 30% when aggregates are broken. Significantly smaller change than expected from Mie ($\times 10$ from $100 \rightarrow 10 \mu m$).

Aggregation in the marine environment

Aggregation is a $[concentration]^2$ phenomena.

Mechanisms for encounter: Brownian motion, differential settling, and turbulent shear.

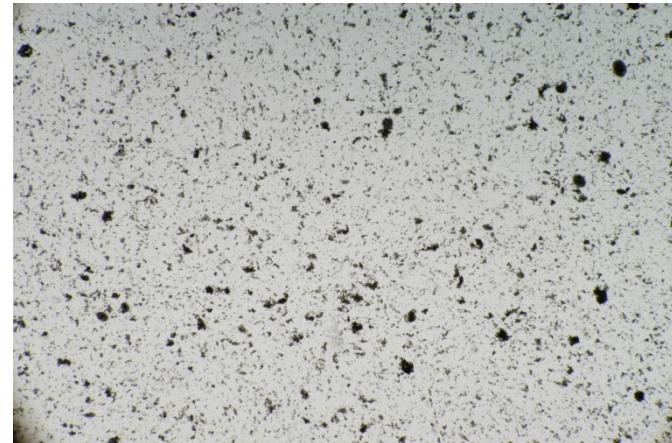
Aggregate sink faster than their component particles.

Aggregates break when shear is too high.

Camera pictures at 1mab at a 12m deep site within 1day:

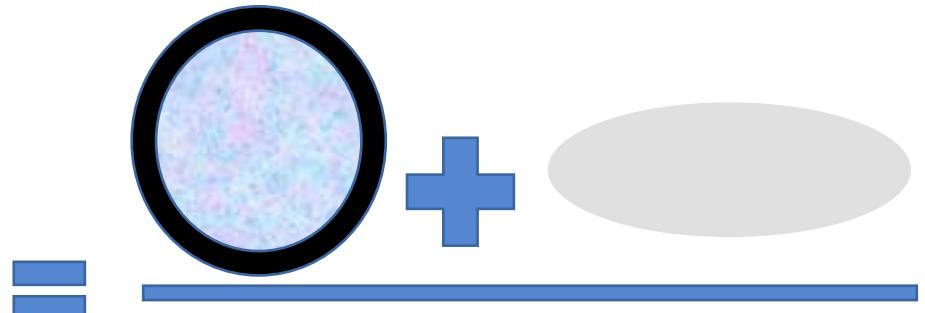


Dominated by $<100\mu\text{m}$ particles



Dominated by $>1000\mu\text{m}$ particles

Aggregate modeling :

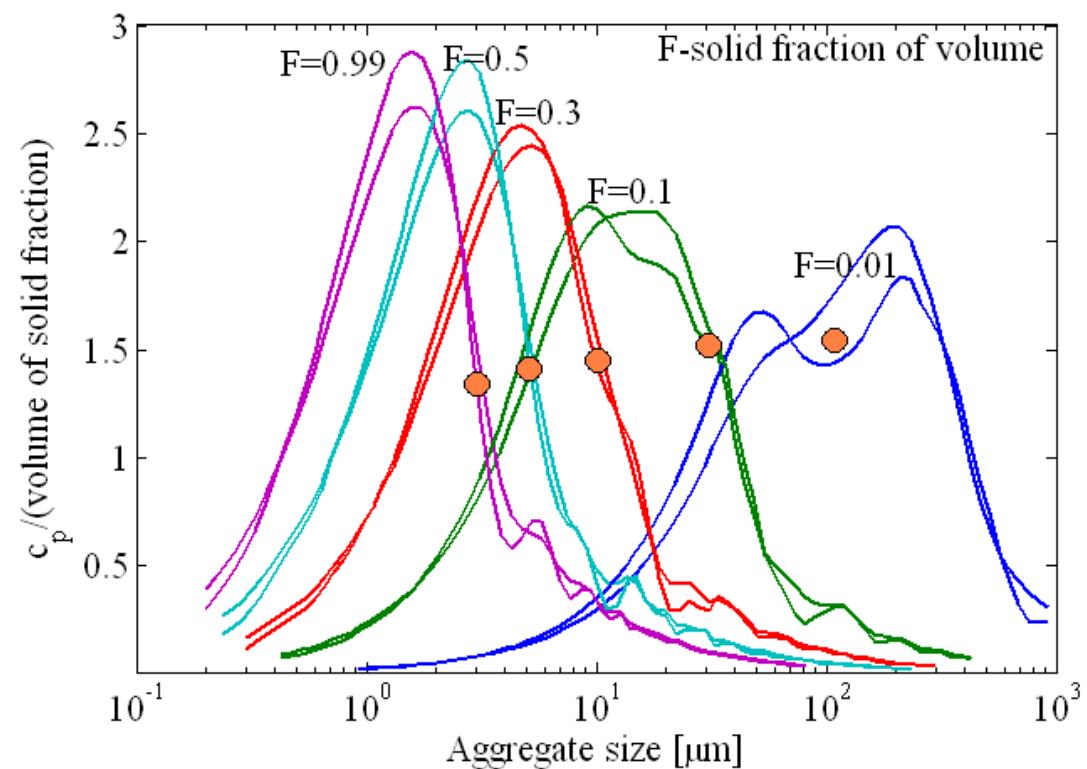


Latimer (1985)

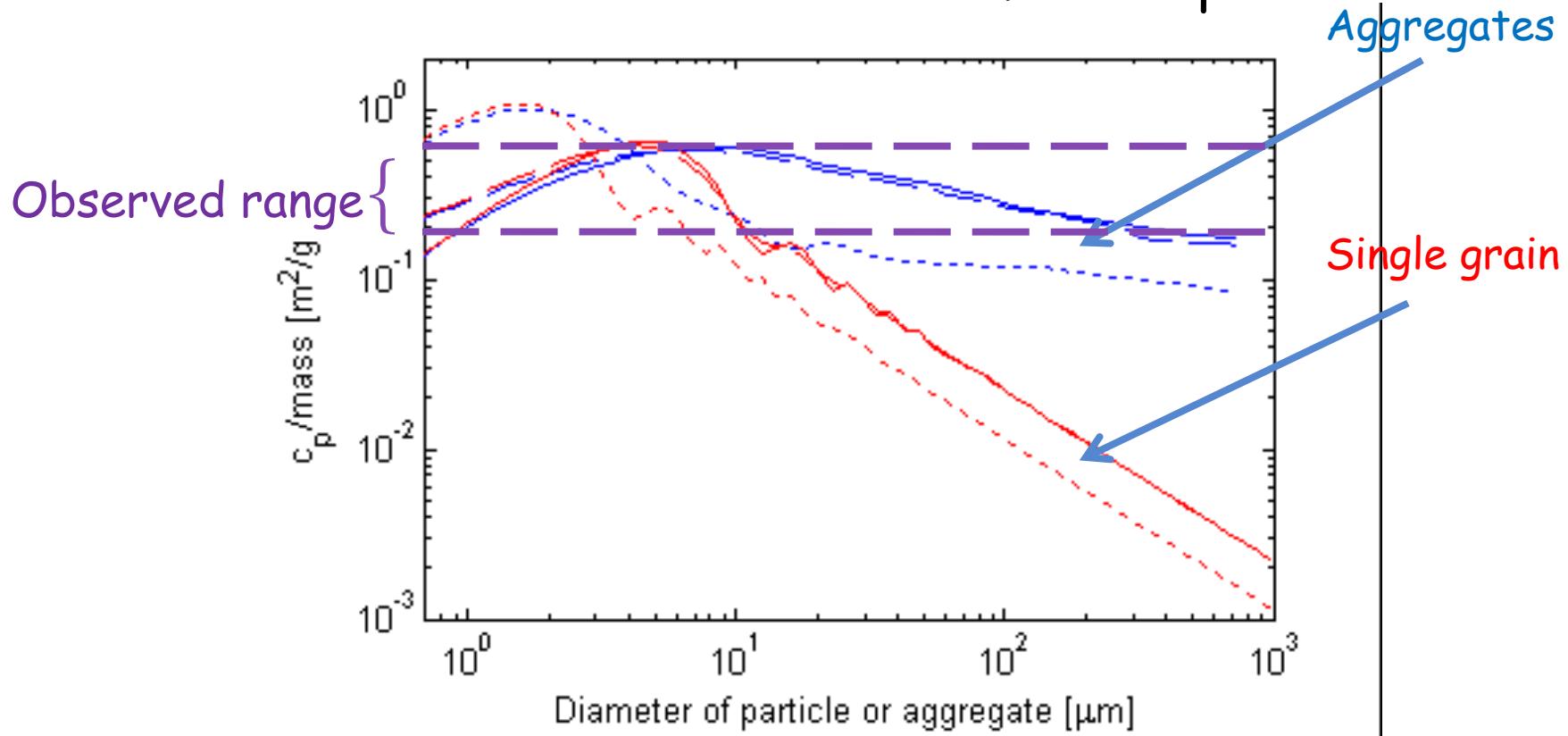
For marine aggregates
size and solid fraction
correlate.



- points having size-F as in Maggi, 2007, or Khelifa and Hill, 2006.

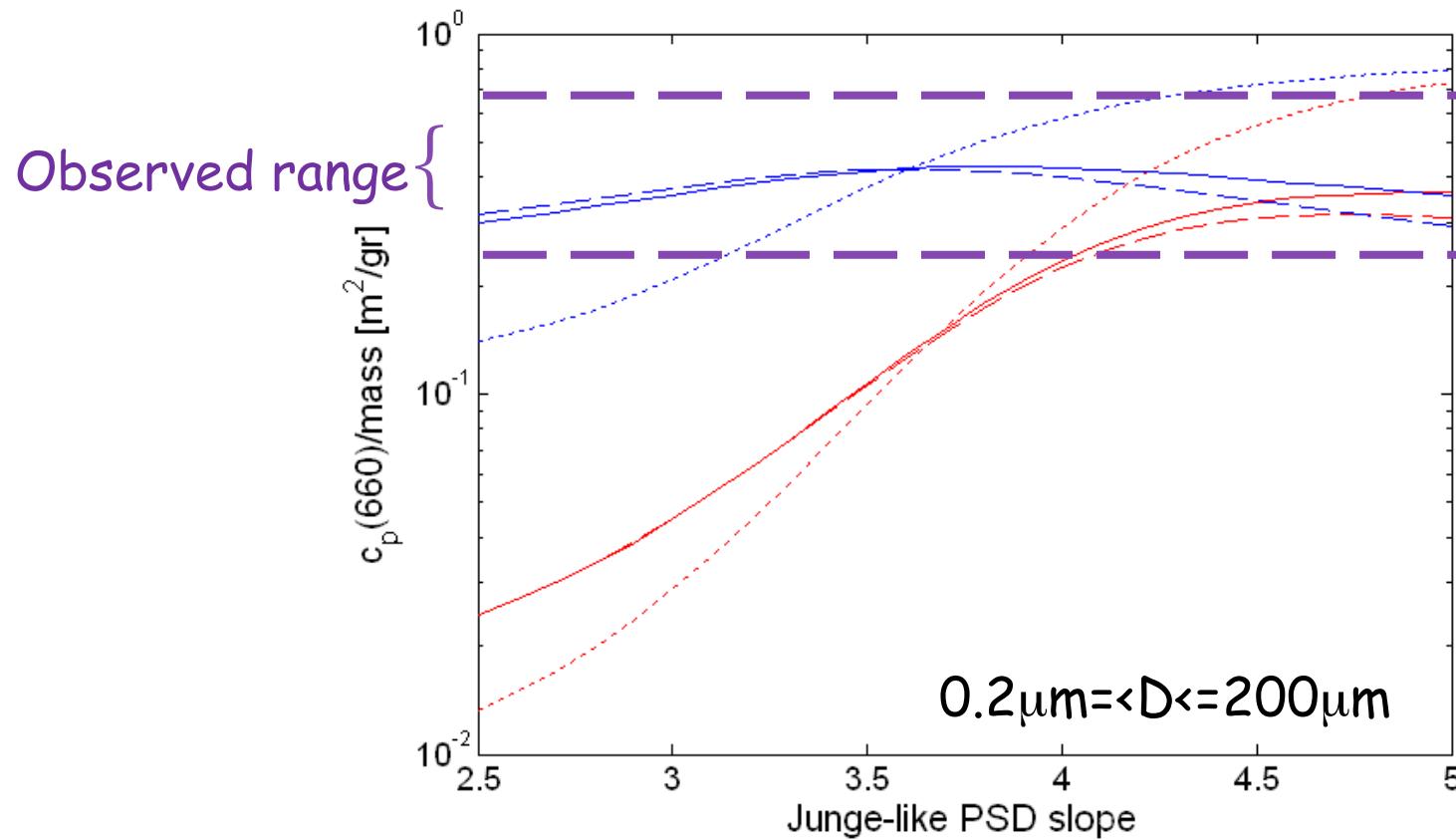


Theoretical calculations: monodispersion



Mass normalized beam attenuation for aggregates assuming a relationship between solid fraction and size as in Khelifa and Hill, 2006 (blue lines) and solid particles (red lines). Solid lines denote particles with $n=1.05+i0.0001$, dashed lines $n=1.05+0.005$ and dotted lines $n=1.15+0.0001$. Each data point represent a population of particle all of a single size.

Theoretical calculations: populations



Mass normalized beam attenuation for populations of aggregates assuming a relationship between solid fraction and size as in Khelifa and Hill, 2006 (blue lines) and populations of solid particles (red lines) both as function of power-law exponent of the disaggregated particle populations. Solid lines denote particles with $n=1.05+i0.0001$, dashed lines $n=1.05+0.005$ and dotted lines $n=1.15+0.0001$.

Note: model is sensitive to size of primary particle, D_{\max} , $F(D_{\max})$ and acceptance angle.

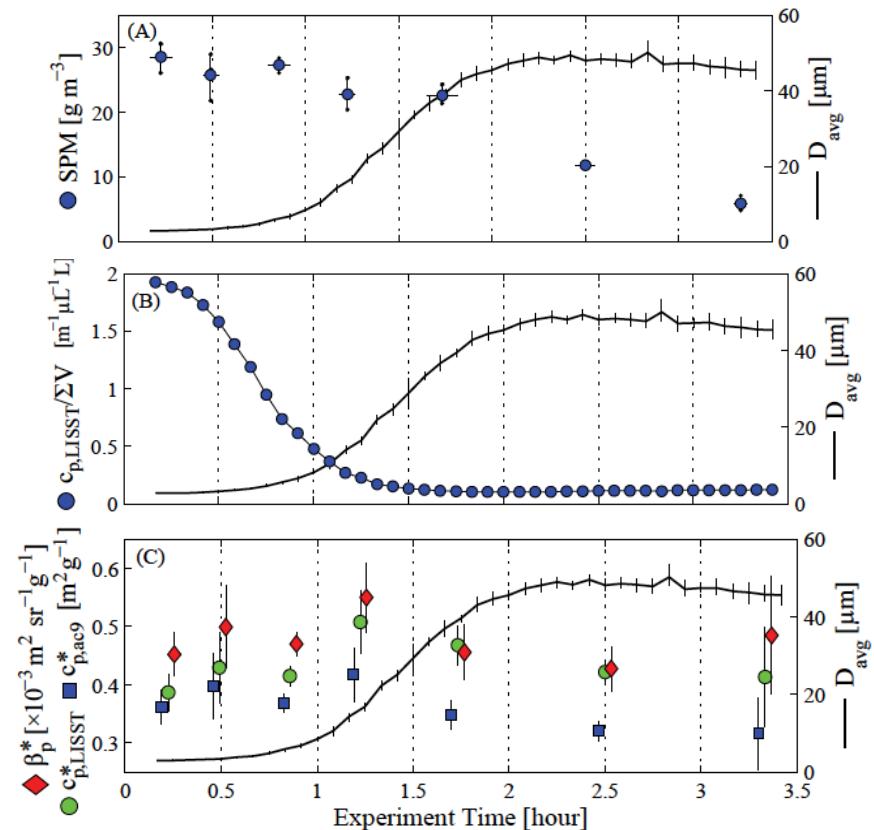
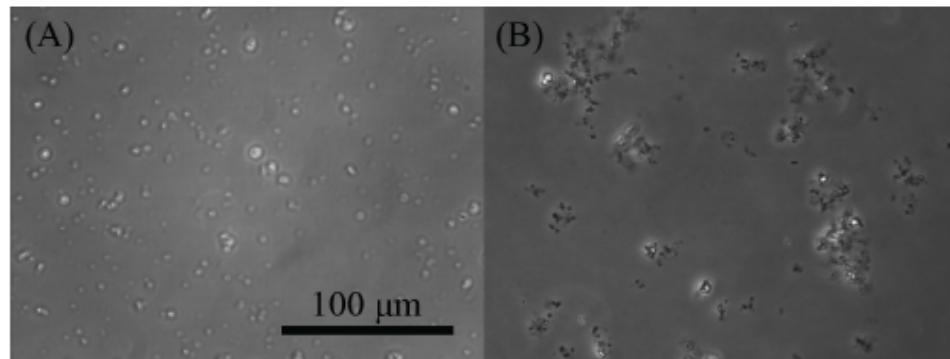
A proxy for aggregate packing

If c_p is a good proxy of mass, and near-forward scattering is a good proxy of volume distribution, than we could obtain an aggregate density proxy by: $c_p/\Sigma \text{volume}$.

Aggregation experiment:

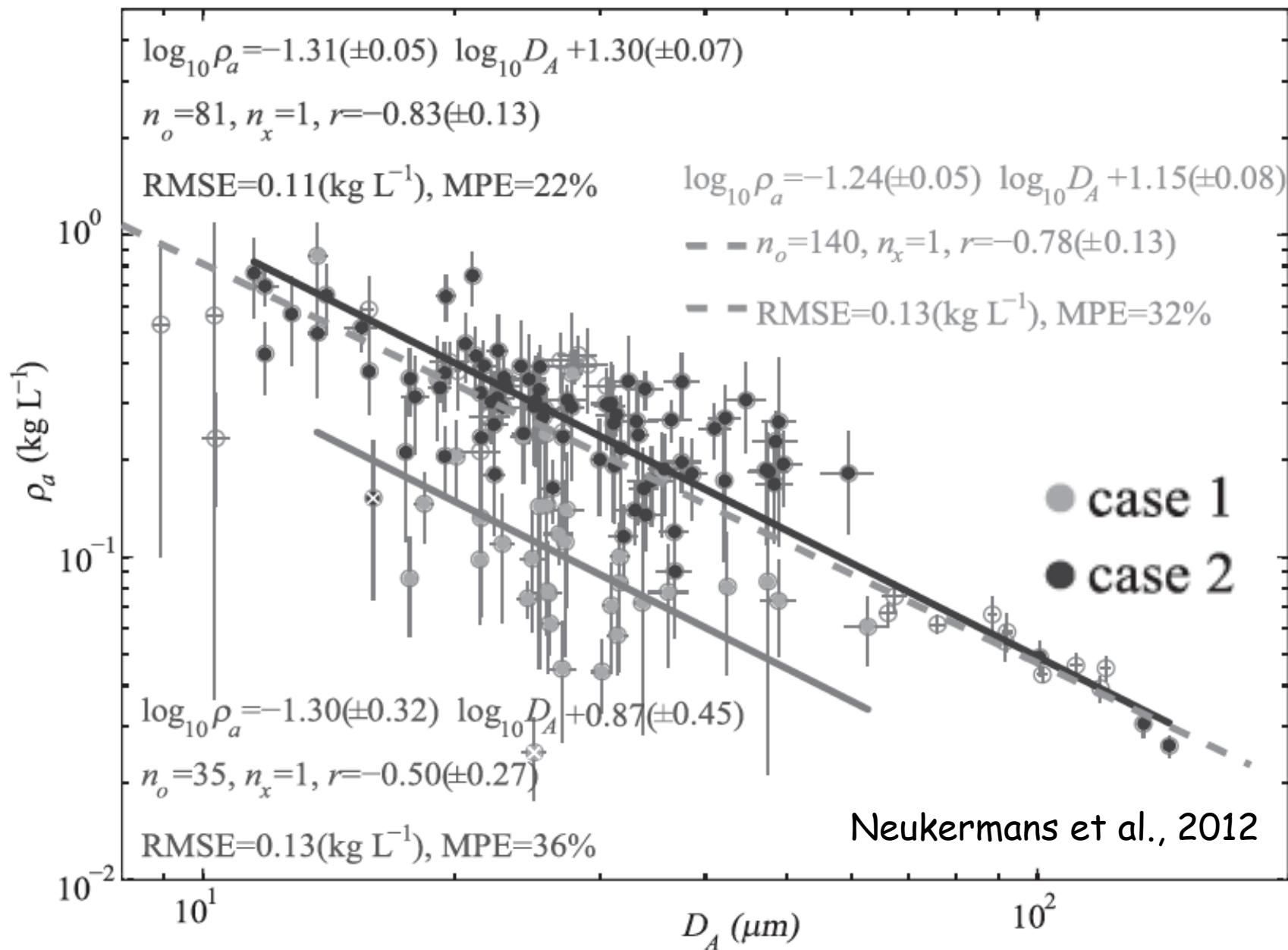
Start with $\langle D \rangle \sim 7 \mu\text{m}$
clay

Add salt



Slade et al., 2011

A proxy for aggregate packing- consistency check



From Inherent Optical Properties to Biogeochemical Properties

Summary of first lecture

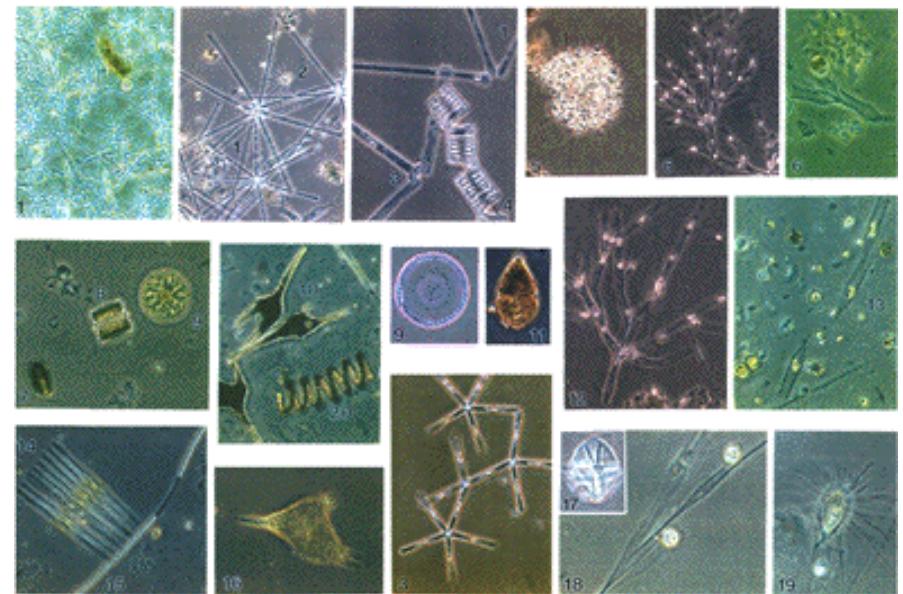
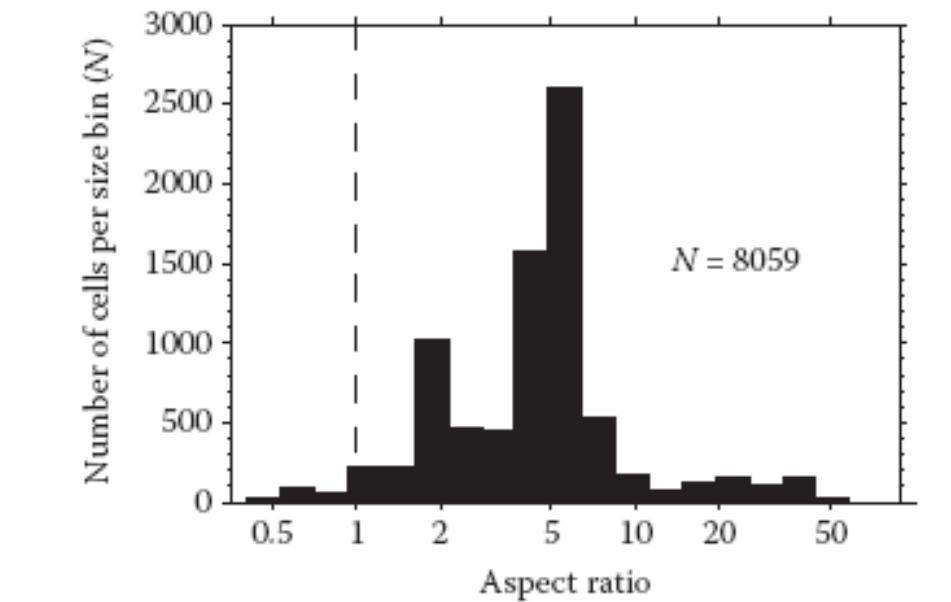
- In this lecture we looked at scattering and attenuation the proxies derived by them.
- Lab studies are critical to test proxies.
- Utility of a proxy is application dependent (tolerance for uncertainties varies).
- Always test the applicability of a proxy before/while you use it.

Questions?

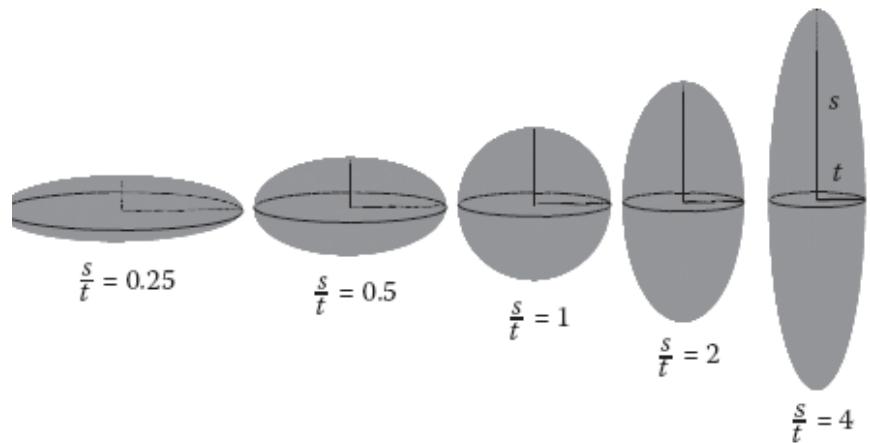
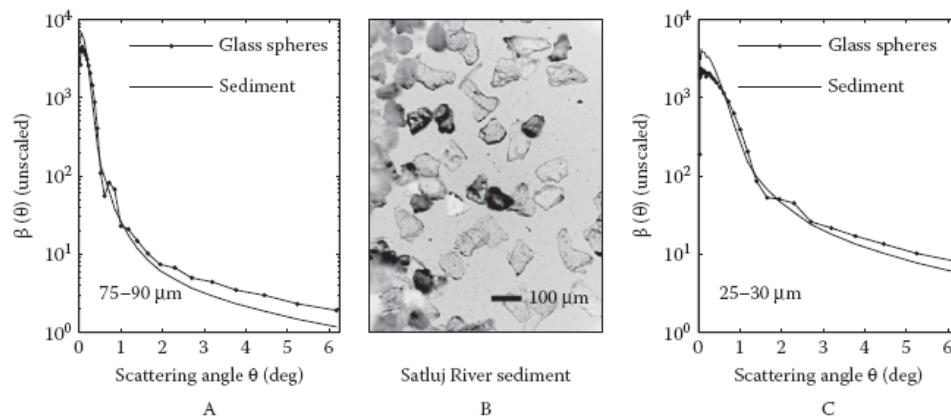
愚者不問，問者不愚。

The fool does not ask, he who asks is no fool

Shape consideration

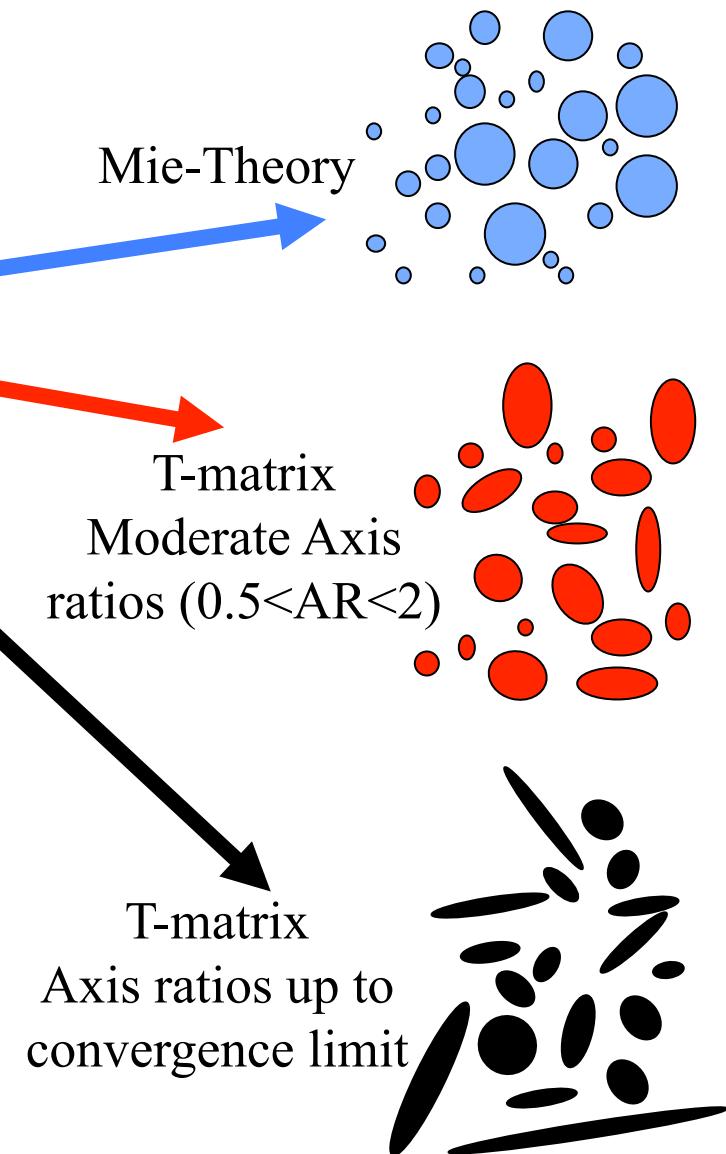
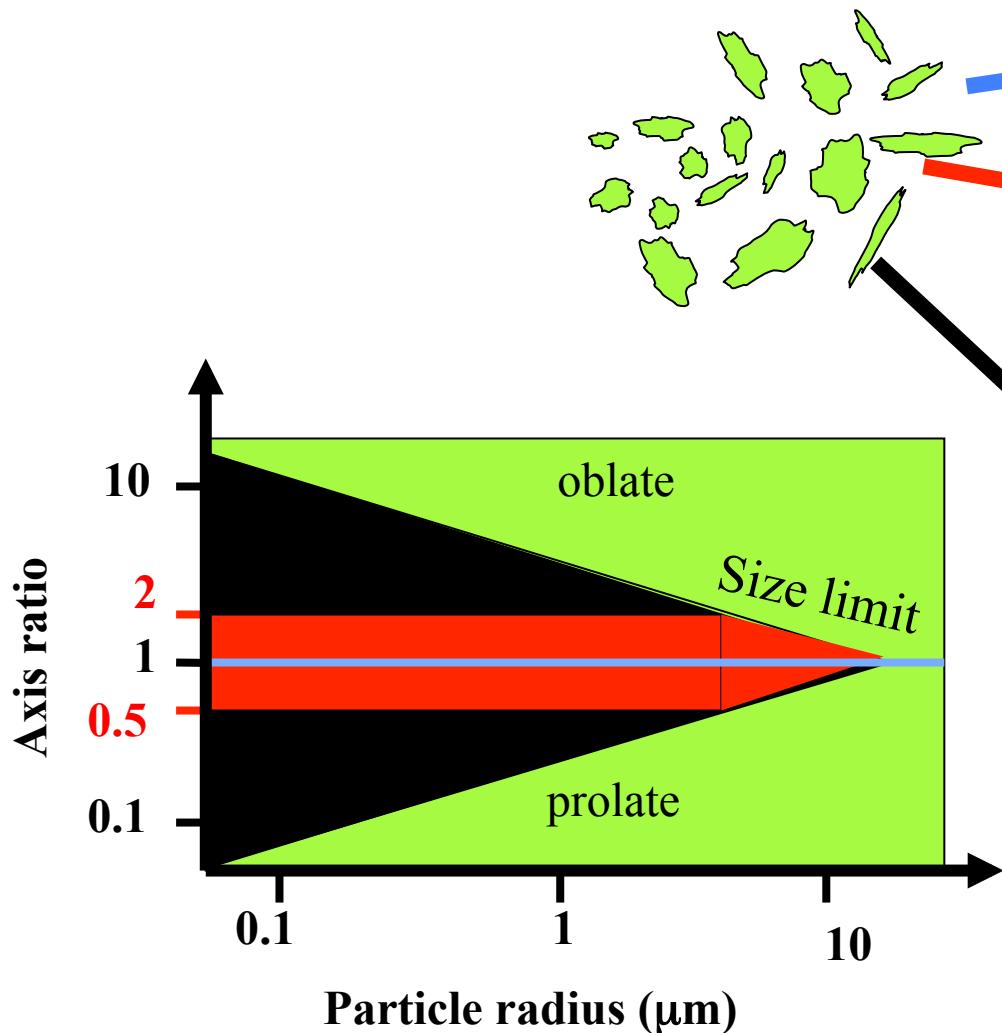


© 1993 Gertrud Cronberg



Clavano et al., 2007

Shape approximations for light scattering calculations



Slide From Volten

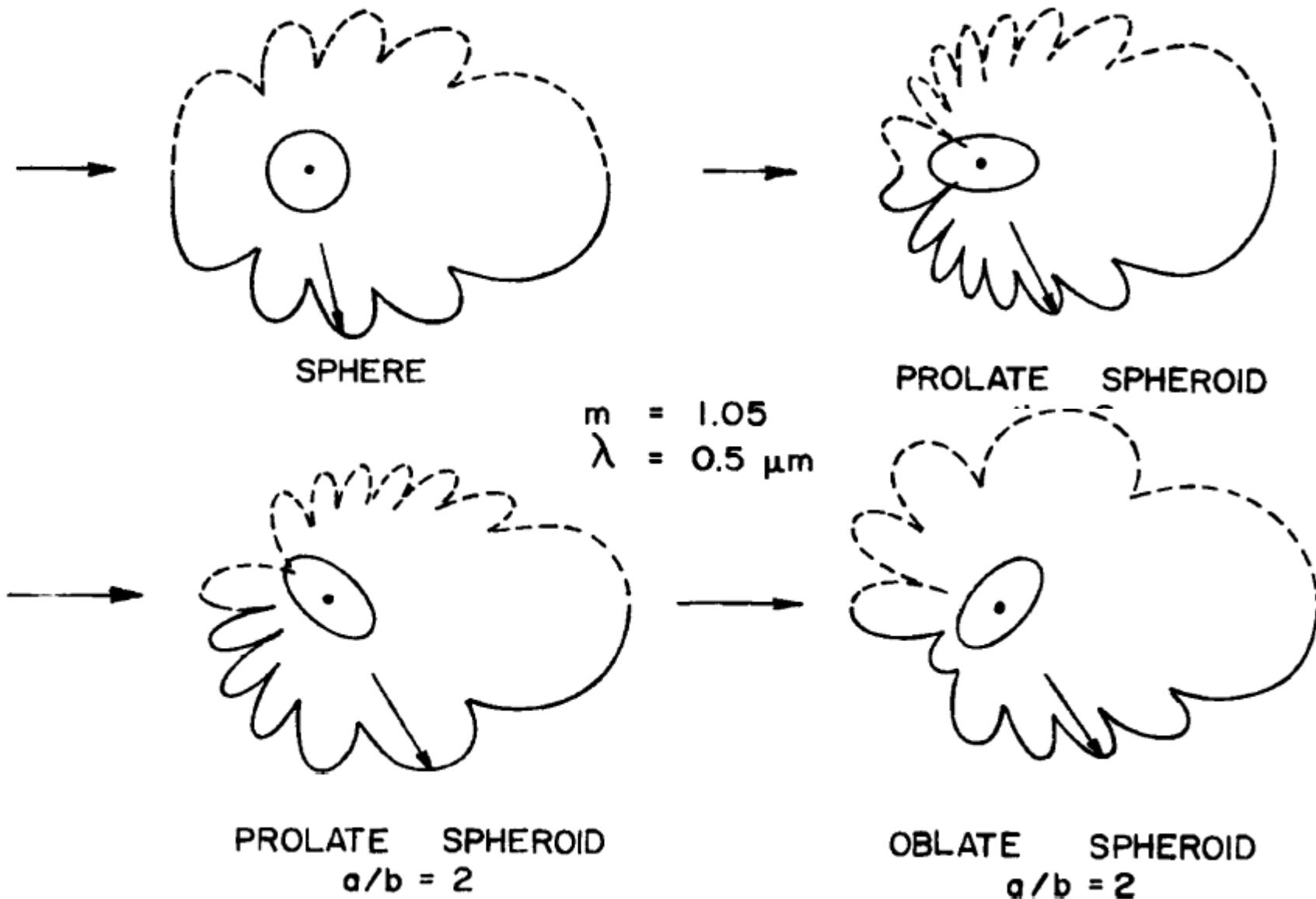
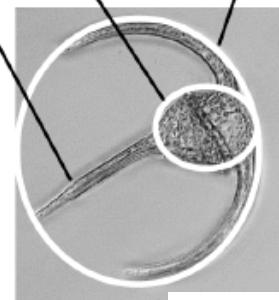
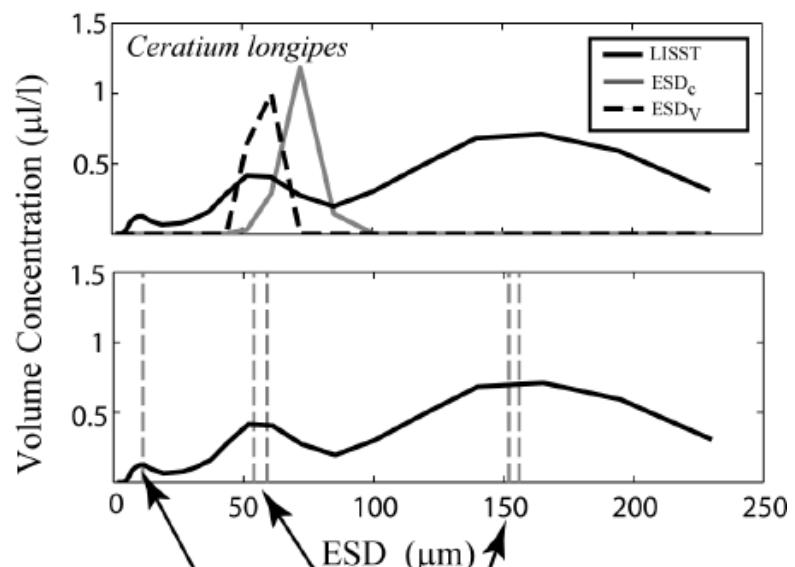
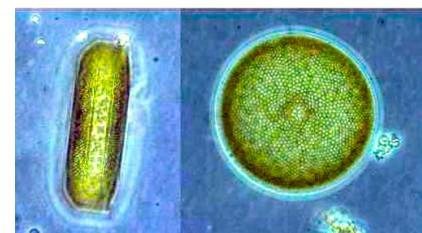
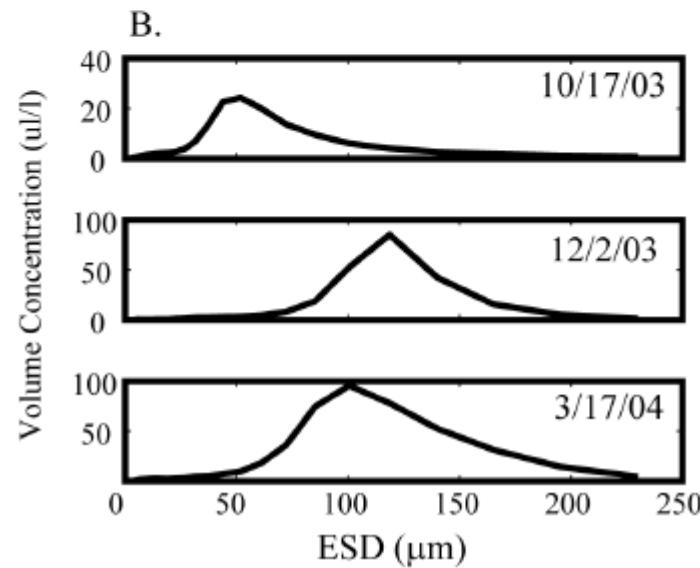
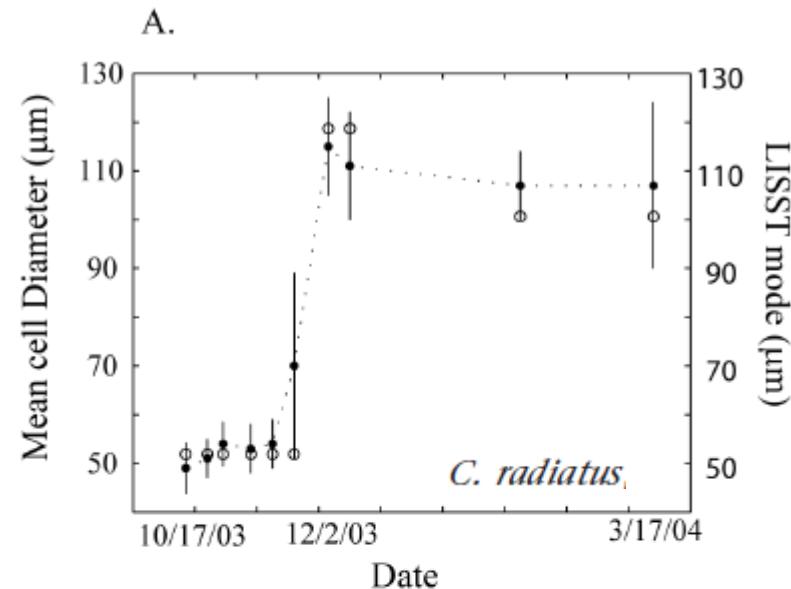
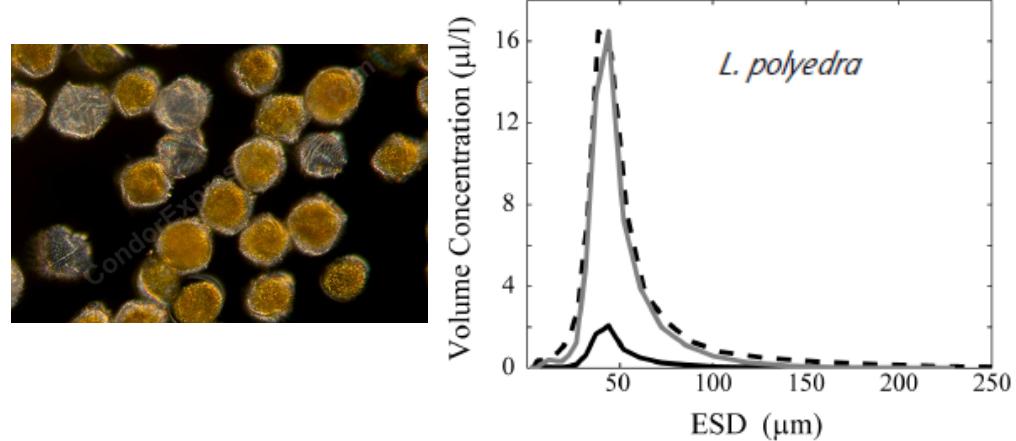


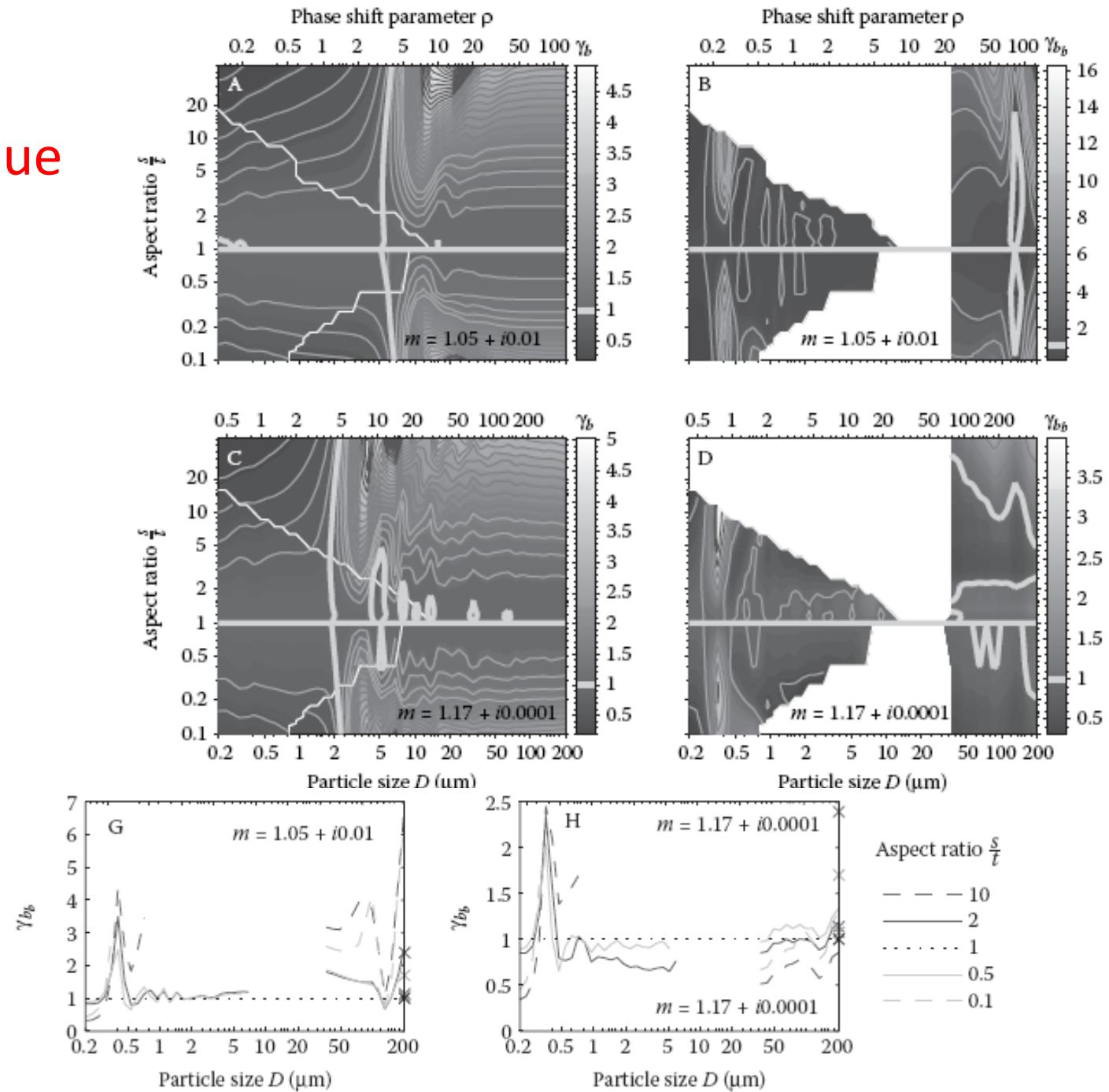
Figure 13.6 Polar scattering diagrams for equal-volume spheroids. The incident light is unpolarized. From Latimer et al. (1978).



Karp-Boss et al., 2007

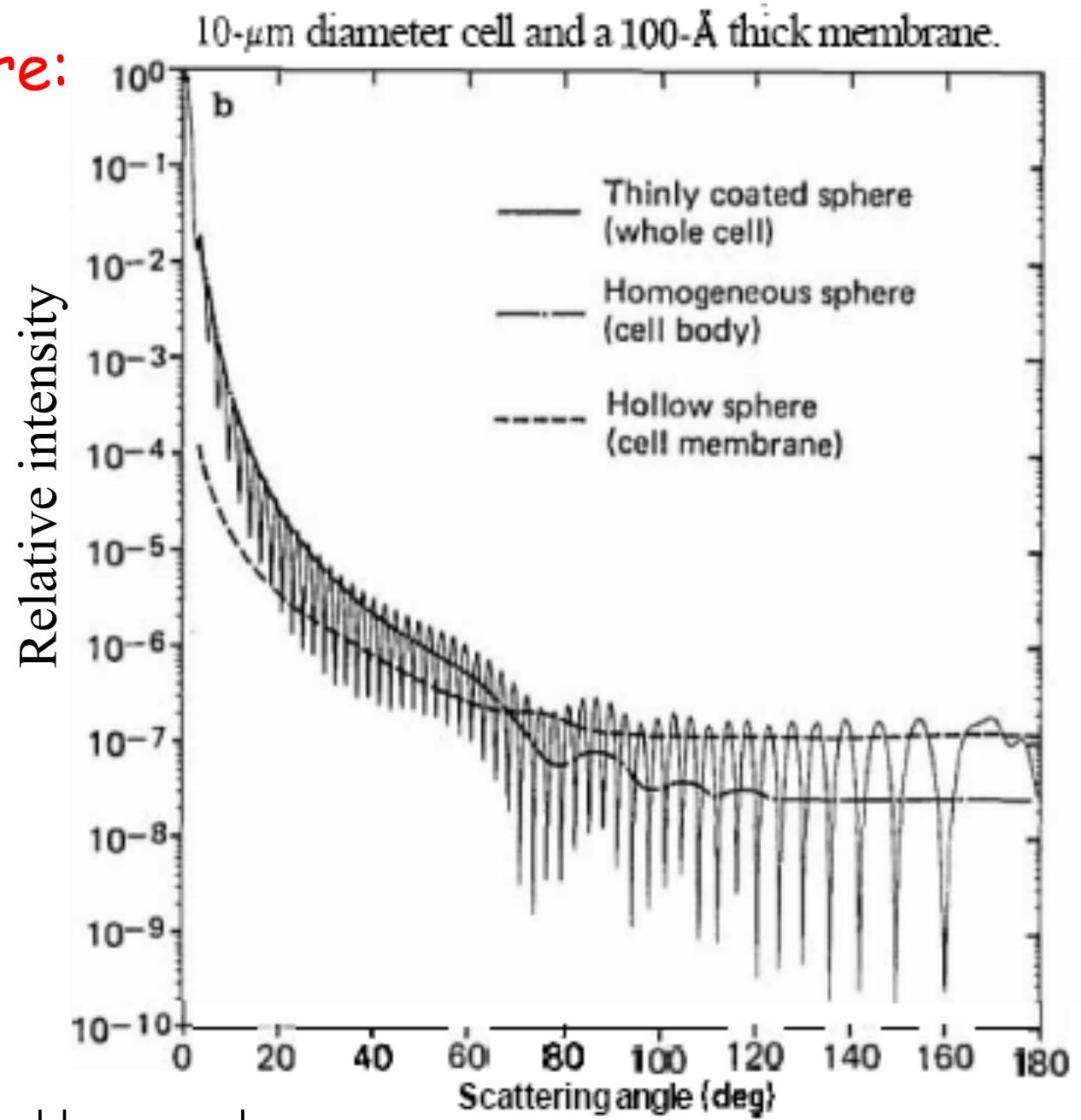


Quantifying differences due to shape:



Clavano et al., 2007

Internal structure:



Backscattering dominated by membrane.

Meyer, 1979